Twisted string algebras and fake flatness

Lennart Schmidt



School of Mathematical and Computer Sciences Heriot-Watt University, Edinburgh

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Self-dual string and fake flatness

- self-dual string equation: $\mathcal{H} = *d\Phi$
- consider 2 models for string Lie 2-algebra:

 $\begin{aligned} & \mathfrak{string}_{\mathrm{sk}} : \mathbb{R} \xrightarrow{0} \mathfrak{g} \ , \\ & \mathfrak{string}_{\Omega} : \widehat{\Omega} \mathfrak{g} \xrightarrow{\mu_1} P_0 \mathfrak{g} \ . \end{aligned}$

- Infinitesimal gauge transformation: $\delta \mathcal{H} = \mu_2(\mathcal{F}, \Lambda)$
- Under categorical equivalence: $\mathcal{H} \to \Phi_0(\mathcal{H}) + \Phi_2(\mathcal{F}, A)$

Fake-flatness condition, $\mathcal{F} = 0$, needed! \leftarrow problematic

- (Higher) Gauge Theory can be encoded in morphisms between differential graded algebras
- ► Morphisms from CE(g) to Ω[•](X) yield potentials and flat curvatures
- ► Gauge transformations are given by homotopies between these

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→ Example:

$$\begin{split} A:\mathsf{CE}(\mathfrak{g})&\longrightarrow\Omega^{\bullet}(X)\\ t^{\alpha}&\longmapsto A^{\alpha}\\ Qt^{\alpha}&=-\tfrac{1}{2}f^{\alpha}_{\beta\gamma}t^{\beta}t^{\gamma}\longmapsto 0=\mathrm{d}A+\tfrac{1}{2}\left[A,A\right] \end{split}$$

Higher Gauge Theory: Weil algebra

Weil algebra $W(\mathfrak{g})$:

- dga: $(\wedge^{\bullet}(\mathfrak{g}^*[1] \oplus \mathfrak{g}^*[2]), Q_{\mathsf{W}})$
- $Q_{\mathsf{W}}|_{\mathfrak{g}^*[1]} = Q_{\mathsf{CE}} + \sigma$ and $Q_{\mathsf{W}}|_{\mathfrak{g}^*[2]} = -\sigma Q_{\mathsf{CE}} \sigma^{-1}$
- $\sigma:\mathfrak{g}^*[1] \to \mathfrak{g}^*[2]$ is the shift isomorphism

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- $\sigma: \mathfrak{g}^*[1] \to \mathfrak{g}^*[2]$ is the shift isomorphism

→ Example:

 \blacktriangleright ordinary Lie algebra $\mathfrak{g},$ generators t^{α} and $r^{\alpha}=\sigma t^{\alpha}$

•
$$Qt^{\alpha} = -\frac{1}{2}f^{\alpha}_{\beta\gamma}t^{\beta} \wedge t^{\gamma} + r^{\alpha}$$

$$\bullet \ Qr^{\alpha} = -f^{\alpha}_{\beta\gamma}t^{\beta} \wedge r^{\gamma}$$

- (Higher) Gauge Theory can be encoded in morphisms between differential graded algebras, i.e. morphisms that commute with Q
- ► Morphisms from W(𝔅) to Ω[•](X) yield potentials, curvatures and Bianchi identities
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➤ Example:

$$\begin{split} (A,\mathcal{F}): \mathbb{W}(\mathfrak{g}) &\longrightarrow \Omega^{\bullet}(X) \\ t^{\alpha}, r^{\alpha} &= \sigma t^{\alpha} &\longmapsto A^{\alpha}, \mathcal{F}^{\alpha} \\ Qt^{\alpha} &= -\frac{1}{2} f^{\alpha}_{\beta\gamma} t^{\beta} t^{\gamma} + r^{\alpha} \longmapsto \mathcal{F}^{\alpha} = \mathrm{d}A + \frac{1}{2} \left[A, A \right] \\ Qr^{\alpha} &= -f^{\alpha}_{\beta\gamma} t^{\beta} r^{\gamma} &\longmapsto (\nabla \mathcal{F})^{\alpha} = 0 \end{split}$$

Higher Gauge Theory: Invariant polynomials

Invariant polynomials inv(g):

- $\blacktriangleright \text{ inv. pol.: } P \in \mathsf{W}(\mathfrak{g})|_{\wedge^{\bullet}(\mathfrak{g}^{*}[2])} \text{ s.t. } Q_{\mathsf{W}}P \in \mathsf{W}(\mathfrak{g})|_{\wedge^{\bullet}(\mathfrak{g}^{*}[2])}$
- $P_1 \sim P_2$ if $P_1 P_2 = Q_W \tau$ for $\tau \in W(\mathfrak{g})|_{\wedge^{\bullet}(\mathfrak{g}^*[2])}$
- ► inv(g) : invariant polynomials mod horizontal equiv.

→ ► identified with characteristic classes

in Chern-Weil theory

▶ sits in short exact sequence:

 $\mathsf{CE}(\mathfrak{g}) \twoheadleftarrow \mathsf{W}(\mathfrak{g}) \hookleftarrow \mathsf{inv}(\mathfrak{g})$

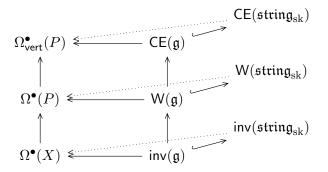
Higher Gauge Theory: g-connection object

Let $P \to X$ be a principal bundle and define g-connection object:

$$\begin{array}{cccc} \Omega^{\bullet}_{\mathsf{vert}}(P) & \xleftarrow{A} & \mathsf{CE}(\mathfrak{g}) \\ & \uparrow & & \uparrow & & \\ \Omega^{\bullet}(P) & \xleftarrow{(A,\mathcal{F})} & \mathsf{W}(\mathfrak{g}) \\ & \uparrow & & \uparrow & & \\ \Omega^{\bullet}(X) & \xleftarrow{<\mathcal{F}>} & \mathsf{inv}(\mathfrak{g}) \end{array} & \xleftarrow{2^{\mathrm{nd}}} & \mathsf{Ehresmann \ condition} \end{array}$$

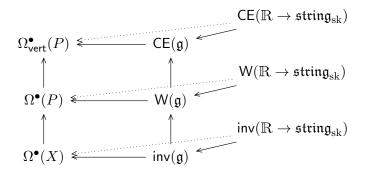
Higher Gauge Theory: Lifting problem

Try to lift to \mathfrak{string}_{sk} :



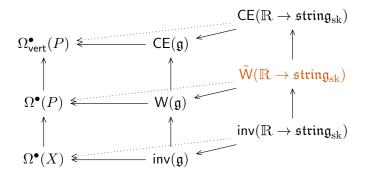
Higher Gauge Theory: Lifting problem

Try to lift to $\mathbb{R} \to \mathfrak{string}_{sk}$:



Higher Gauge Theory: Lifting problem

Try to lift to $\mathbb{R} \to \mathfrak{string}_{sk}$:



Modified string algebra: Skeletal model

This yields:

$$\widetilde{\mathfrak{string}}_{sk} = \ \mathbb{R}[2] \longrightarrow \mathbb{R}[1] \overset{0}{\longrightarrow} \mathfrak{g}$$
 ,

together with:

 $\mathcal{F} = dA + \frac{1}{2}\mu_2(A, A) ,$ $\mathcal{H} = dB + \frac{1}{6}\mu_3(A, A, A) - \kappa(A, \mathcal{F}) + \mu_1(C) ,$ $\mathcal{G} = dC ,$ $d\mathcal{F} = -\mu_2(A, \mathcal{F}) ,$ $d\mathcal{H} = -\kappa(\mathcal{F}, \mathcal{F}) + \mu_1(\mathcal{G}) ,$ $d\mathcal{G} = 0 .$

 \implies SDS equation gauge covariant even for non-vanishing ${\cal F}$.

Modified string algebra: Loop model

Analogous story yields:

$$\widetilde{\mathfrak{string}}_{\Omega} = \mathbb{R}[2] \longrightarrow \hat{\Omega}(\mathfrak{g})[1] \xrightarrow{\mu_1} P_0 \mathfrak{g} ,$$

together with:

$$\mathcal{F} = dA + \frac{1}{2}\mu_2(A, A) + \mu_1(B) ,$$

$$\mathcal{H} = dB + \mu_2(A, B) - \kappa(A, \mathcal{F}) + \mu_1(C) ,$$

$$\mathcal{G} = dC ,$$

$$d\mathcal{F} = -\mu_2(A, \mathcal{F}) + \mu_1(\mathcal{H}) ,$$

$$d\mathcal{H} = -\kappa(\mathcal{F}, \mathcal{F}) + \mu_1(\mathcal{G}) ,$$

$$d\mathcal{G} = 0 .$$

Consider:



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