Global Solutions to Stochastic Reaction-Diffusion Equations

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In this talk, I will present some recent results on the global existence of solutions to stochastic reaction-diffusion(/stochastic heat) equations with super-linear drift and multiplicative noise.

This talk is based on the joint work with Robert Dalang and Davar Khoshnevisan.

Let ξ denote the space-time white noise on $\mathbb{R}_+ \times [0, 1]$, and consider the parabolic stochastic partial differential equation

$$\dot{u}(t,x) = \frac{1}{2}u''(t,x) + b(u(t,x)) + \sigma(u(t,x))\xi(t,x), \quad (1)$$

t > 0, $x \in (0, 1)$, subject to the homogeneous Dirichlet boundary condition,

$$u(t, 0) = u(t, 1) = 0$$
 for all $t > 0$,

and the initial condition $u(0,x) = u_0(x)$, $x \in [0,1]$. Throughout, $\sigma : \mathbb{R} \to \mathbb{R}$ is assumed to be a nonrandom and measurable function, and $b : \mathbb{R} \to \mathbb{R}$ is assumed to be nonrandom and measurable. We assume further that the initial function $u_0 : [0,1] \to \mathbb{R}$ is always nonrandom. It is well known that if, in addition, b, σ have at most linear growth—that is if $|b(z)|, |\sigma(z)| = O(|z|)$ as $|z| \to \infty$ —then any local solution of (1) is a global one. A few years ago, Bonder and Groisman proved the following interesting complement.

Theorem. [Bonder and Groisman] Suppose, in addition, that σ is a nonzero constant, b is a nonnegative convex function, and satisfies either $\int_{1}^{\infty} dz/b(z) < \infty$ or $\int_{-\infty}^{-1} dz/b(z) < \infty$, or both, and the initial function u_0 is nonnegative, continuous on [0, 1], and vanishes on $\{0, 1\}$. Then there exists an almost surely finite random time τ such that

$$\int_0^1 |u(t,x)|^2 \,\mathrm{d} x = \infty$$
 for every $t > au$.

Our goal is to prove that the preceding result of Bonder–Groisman is in a certain sense optimal. In fact, we introduce two rather different methods which show that, under two different sets of natural conditions, if $|b(z)| = O(|z| \log |z|)$ then the solution to (1) does not blow up at finite time. Let us first recall the definition of a \mathbb{L}^2 -solution. **Definition 1**. Let τ be a stopping time. A $L^2[0, 1]$ -valued continuous, adapted random field $\{u(t, \cdot), t \in [0, \tau)\}$ is called a solution to equation (1) if for every test function $\phi \in C_0^2(0, 1)$,

$$\int_{0}^{1} u(t,x)\phi(x)dx = \int_{0}^{1} u_{0}(x)\phi(x)dx + \frac{1}{2}\int_{0}^{t}\int_{0}^{1} u(s,x)\phi''(x)dx + \int_{0}^{t}\int_{0}^{1} b(u(s,x))\phi(x)dx + \int_{0}^{t}\int_{0}^{1} \sigma(u(s,x))\phi(x)\xi(dsdx)$$
(2)

a.s. for all $t \in [0, \tau)$.

Our first result is stated as follows.

Theorem 1. Suppose that $u_0 \in L^2[0,1]$, $\sigma : \mathbb{R} \to \mathbb{R}$ is bounded, and $|b(z)| = O(|z| \log |z|)$ as $|z| \to \infty$. Then, any \mathbb{L}^2 -solution u of (1) does not blow up in finite time. **Sketch of the proof.** We will appeal to the logarithmic Sobolev inequality of Gross [2] in the following form: For every $\varepsilon \in (0, 1)$ and infinitely-differentiable functions $h : [0, 1] \to \mathbb{R}$ that vanish continuously on $\{0, 1\}$,

$$\int_0^1 |h(x)|^2 \log |h(x)| \,\mathrm{d} x \leqslant \varepsilon \|h'\|_{\mathbb{L}^2}^2 + \frac{1}{4} \log(1/\varepsilon) \|h\|_{\mathbb{L}^2}^2 + \|h\|_{\mathbb{L}^2}^2 \log\left(\|h\|_{\mathbb{L}^2}^2\right),$$

where $0 \log 0 := 0$.

For every R > 0, consider the stopping times,

$$au(R):=\inf\left\{t>0: \|u(t)\|_{\mathbb{L}^2}>R
ight\} \quad ext{and let} \quad au:=\lim_{R o\infty} au(R).$$

Our goal is to prove that $P\{\tau = \infty\} = 1$. We will do this by proving that $P\{\tau < T\} = 0$ for every positive constant T.

For every constant R > 0 consider the following stochastic PDE with random forcing and no reaction term:

$$\dot{v}_{R}(t,x) = \frac{1}{2}v_{R}''(t,x) + \sigma(u(t \wedge \tau(R), x))\xi(t,x), \quad [0 < t < \tau, \ 0 \le x \le 1].$$
(3)

The solution process $t \mapsto v_R(t)$ satisfies the following random integral equation:

$$v_{\mathcal{R}}(t,x) = \int_{(0,t)\times[0,1]} G_{t-s}(x,y)\sigma(u(s\wedge\tau(\mathcal{R}),y))\xi(\mathrm{d} s\,\mathrm{d} y). \quad (4)$$

Here, the function $G: (0, \infty) \times [0, 1]^2 \to \mathbb{R}_+$ denotes the heat kernel. We will use $\mathcal{G} := \{\mathcal{G}_t\}_{t \ge 0}$ to denote the corresponding heat semigroup.

Define, for every fixed R > 0,

$$d_R := u - v_R.$$

We may observe that d_R solves the following random heat equation: For all $t \in [0, \tau)$,

$$\dot{d}_{R}(t) = \frac{1}{2} d_{R}''(t) + b \left(v_{R}(t) + d_{R}(t) \right),$$
(5)

subject to $d_R(0) = u_0$. Next we consider the Lyapunov function,

$$\Phi(r) := \exp\left(\int_0^r \frac{\mathrm{d}z}{1+z\log_+ z}\right),\,$$

defined for every r > 0.

Sketch of the proof

Choose and fix some T > 0. Since σ is a bounded measurable function, we can show that

$$A := \sup_{R>0} \mathbb{E} \left(\sup_{t \in [0,T]} \sup_{x \in [0,1]} |v_R(t,x)| \right) < \infty.$$
 (6)

Consider the stopping time

$$au_M(R):=\inf\left\{t>0: \sup_{x\in[0,1]}|v_R(t\,,x)|>M
ight\} \qquad ext{for every } M>0.$$

It follows from (6) and the Chebyshev inequality that

$$\sup_{R>0} \operatorname{P}\left\{\tau_M(R) < T\right\} \leqslant \frac{A}{M}.$$
(7)

Temporarily define two random space-time functions D and V as $D(t) := d_R (t \wedge \tau(R) \wedge \tau_M(R)), \quad V(t) := v_R (t \wedge \tau(R) \wedge \tau_M(R))$ for $0 \le t \le T$, all the time suppressing the dependence of D and V on (R, M), as well as on the spatial variable $x \in [0, 1]$.

Sketch of the proof

We are able to justify the use of the chain rule to get that for every $t \in [0, T]$,

$$\|D(t)\|_{\mathbb{L}^{2}}^{2} = \|u_{0}\|_{\mathbb{L}^{2}}^{2} - 2\int_{0}^{t} \|D'(s)\|_{\mathbb{L}^{2}}^{2} ds$$

+2 $\int_{0}^{t} \langle b(V(s) + D(s)), D(s) \rangle_{\mathbb{L}^{2}} ds.$ (8)

A second application of the chain rule yields

$$\Phi\left(\|D(t)\|_{\mathbb{L}^{2}}^{2}\right) = \Phi\left(\|u_{0}\|_{\mathbb{L}^{2}}^{2}\right) - 2\int_{0}^{t}\Phi'\left(\|D(s)\|_{\mathbb{L}^{2}}^{2}\right)\left\|D'(s)\right\|_{\mathbb{L}^{2}}^{2} \mathrm{d}s$$
$$+ 2\int_{0}^{t}\Phi'\left(\|D(s)\|_{\mathbb{L}^{2}}^{2}\right)\left\langle b(V(s) + D(s)), D(s)\right\rangle_{\mathbb{L}^{2}} \mathrm{d}s.$$
(9)

Using the growth condition of the drift b we can show that

$$\langle b(V(s) + D(s)), D(s) \rangle_{\mathbb{L}^2} \leqslant \overline{C} \left\{ \|D(s)\|_{\mathbb{L}^2 \log \mathbb{L}}^2 + \|D(s)\|_{\mathbb{L}^2}^2 + 1 \right\},$$

uniformly for all $s \in [0, T]$, where \overline{C} is a non-random and finite constant, and depends only on (C_b, M) . Thus, we may apply the logarithmic Sobolev inequality to get

$$egin{aligned} &\langle b(V(s)+D(s)),D(s)
angle_{\mathbb{L}^2}\leqslant \|D'(s)\|_{\mathbb{L}^2}^2\ &+c_*\left\{\|D(s)\|_{\mathbb{L}^2}^2+\|D(s)\|_{\mathbb{L}^2}^2\log_+\left(\|D(s)\|_{\mathbb{L}^2}^2
ight)+1
ight\}, \end{aligned}$$

uniformly for all $s \in [0, T]$, where c_* is a non-random and finite constant, and depends only on (C_b, M) .

Sketch of the proof

We can deduce the following from (9):

$$\Phi\left(\|D(t)\|_{\mathbb{L}^{2}}^{2}\right) \leqslant \Phi\left(\|u_{0}\|_{\mathbb{L}^{2}}^{2}\right) + C \int_{0}^{t} \Phi'\left(\|D(s)\|_{\mathbb{L}^{2}}^{2}\right) \\ \times \left\{1 + \|D(s)\|_{\mathbb{L}^{2}}^{2} \log_{+}\left(\|D(s)\|_{\mathbb{L}^{2}}^{2}\right)\right\} \mathrm{d}s.$$
(10)

But $\Phi'(r)[1 + r \log_+ r] = \Phi(r)$ for all $r \ge 0$. Therefore, the preceding inequality implies that

$$\Phi\left(\|D(t)\|_{\mathbb{L}^2}^2\right) \leqslant \Phi\left(\|u_0\|_{\mathbb{L}^2}^2\right) + C\int_0^t \Phi\left(\|D(s)\|_{\mathbb{L}^2}^2\right) \mathrm{d}s,$$

uniformly for all $t \in [0, T]$, where the implied constant is non-random and finite, and depends only on (C_b, M, T) . It follows from Gronwall's inequality, that $\sup_{t \in [0,T]} \Phi(||D(t)||_{\mathbb{L}^2}^2)$ is a.s. bounded from above by a non-random finite number $B(C_b, M, T)$, that depends only on (C_b, M) , whence

$$\sup_{R>0} \mathbb{E}\left[\Phi\left(\|d(T \wedge \tau(R) \wedge \tau_M(R)\|_{\mathbb{L}^2}^2\right)\right] \leqslant B(C_b, M, T).$$
(11)

On the other hand, we can show that

 $\|d(T \wedge \tau(R) \wedge \tau_M(R))\|_{\mathbb{L}^2} \ge R - M$

a.s. on the event $\{\tau(R) \leqslant T \leqslant \tau_M(R)\}$, whence

$$\Phi\left(\left\|d(T\wedge au(R)\wedge au_{M}(R))
ight)\|_{\mathbb{L}^{2}}^{2}
ight)\geqslant\Phi\left((R-M)^{2}
ight)$$

a.s. on $\{\tau(R) \leq T \leq \tau_M(R)\}$ as long as R > M.

Combine this with (11) to see that

$$P\left\{\tau(R) \leqslant T \leqslant \tau_M(R)\right\} \leqslant \frac{B(C_b, M, T)}{\Phi\left((R-M)^2\right)} \quad \text{for all } R > M > 0.$$

The preceding inequality and (7) together show that

$$\mathrm{P}\left\{\tau(R)\leqslant T\right\}\leqslant \frac{B(C_b,M,T)}{\Phi\left((R-M)^2\right)}+\frac{A}{M}$$

for all R > M. We first let $R \to \infty$ and then $M \to \infty$ in order to see that

$$\mathrm{P}\{\tau < T\} = \lim_{R \to \infty} \mathrm{P}\{\tau(R) < T\} = 0.$$

Since T > 0 is arbitrary, this proves the theorem.

In this second part, I will introduce another main result in the L^{∞} -setting under a different set of conditions. The approach is also quite different. First let us recall the definition of the solution. **Definition 2.** A random field solution to (1) is a jointly measurable and adapted space-time process $u := \{u(t, x)\}_{(t,x) \in \mathbb{R}_+ \times [0,1]}$ such that, for all $(t, x) \in \mathbb{R}_+ \times [0,1]$,

$$\begin{split} u(t,x) &= (\mathcal{G}_t u_0)(x) + \int_{(0,t)\times(0,1)} G_{t-s}(x,y) b(u(s,y)) \, \mathrm{d}s \, \mathrm{d}y \\ &+ \int_{(0,t)\times(0,1)} G_{t-s}(x,y) \sigma(u(s,y)) \, \xi(\mathrm{d}s \, \mathrm{d}y), \end{split}$$

almost surely, where $\{\mathcal{G}_t\}_{t\geq 0}$ and G are respectively the heat semigroup and heat kernel for the Dirichlet Laplacian.

For every globally Lipschitz function $f : \mathbb{R} \to \mathbb{R}$, there are constants c(f) and L(f) such that

$$|f(z)| \leqslant c(f) + L(f)|z|,$$
 for all $z \in \mathbb{R}$. (12)

One possibility is to take c(f) = |f(0)| and L(f) = Lip(f), but often, L(f) can be chosen strictly smaller than Lip(f)

Here is our second main result: **Theorem 2**. Suppose that:

- ► $u_0 \in \cup_{0 < \alpha \leqslant 1} \mathbb{C}_0^{\alpha};$
- ▶ b and σ are locally Lipschitz functions such that $|b(z)| = O(|z| \log |z|)$ as $|z| \to \infty$; and
- ► $|\sigma(z)| = o\left(|z|(\log |z|)^{1/4}\right)$ as $|z| \to \infty$.

Then, the SPDE (1) has a unique random field solution in $L^{\infty}[0,1]$. In fact, *u* has a continuous modification that satisfies

$$\sup_{t\in[0,T]}\sup_{x\in[0,1]}|u(t,x)|<\infty\qquad\text{a.s., for all }T\in(0,\infty). \tag{13}$$

To prove Theorem 2, we first establish some rather precise moment bounds of the solution of the stochastic heat equation in the classical case that

b and σ are globally Lipschitz continuous.

We consider only the case that

$$L(b) \ge 4L(\sigma)^4 > 0. \tag{14}$$

Proposition 3. The following logical implication is valid:

$$u_0 \in \bigcup_{0 < \alpha \leqslant 1} \mathbb{C}_0^{lpha} \implies \operatorname{P}\left\{u(t) \in \bigcup_{0 < lpha \leqslant 1} \mathbb{C}_0^{lpha} ext{ for all } t > 0
ight\} = 1.$$

Set

$$egin{aligned} \mathcal{M}_1 &:= c(b) + c(\sigma); \quad \mathcal{M}_2 &:= L(b) + L(\sigma); ext{ and} \ \mathcal{M}_3 &:= \|u_0\|_{\mathbb{L}^\infty} + rac{c(b)}{L(b)} + rac{c(\sigma)}{L(\sigma)}. \end{aligned}$$

Proposition 4. Choose and fix $\alpha \in (0, 1]$. There exists a finite universal constant *A*—independent of (b, σ) —such that

$$\sup_{0 \leq x < x' \leq 1} \operatorname{E} \left(\left| \frac{u(t, x) - u(t, x')}{|x' - x|^{\alpha \wedge (1/2)}} \right|^k \right)$$

$$\leq A^k \left(\|u_0\|_{\mathbb{C}_0^{\alpha}}^k + k^{k/2} \mathcal{M}_1^k + k^{k/2} \mathcal{M}_2^k \mathcal{M}_3^k e^{AkL(b)t} \right),$$
(15)

uniformly for all $u_0 \in \mathbb{C}_0^{\alpha}$, $t \ge 0$, and $k \in [2, \sqrt{L(b)}/L(\sigma)^2]$.

Proposition 5. Fix $T_0 > 0$. Choose and fix some $\alpha \in (0, 1]$, and define $\mu := \min(\frac{1}{4}, \frac{1}{2}\alpha)$. Then there exists a finite constant *A*—independent of (b, σ) —such that

$$\sup_{x \in [0,1]} \operatorname{E} \left(\left| \frac{u(T,x) - u(t,x)}{(T-t)^{\mu}} \right|^{k} \right) \\ \leqslant A^{k} \left(\|u_{0}\|_{\mathbb{C}_{0}^{\alpha}}^{k} + k^{k/2} \left[\mathcal{M}_{1}^{k} + \mathcal{M}_{2}^{k} \mathcal{M}_{3}^{k} \mathrm{e}^{\mathcal{A}kL(b)(T)} \right] \right), \quad (16)$$

for all $u_0 \in \mathbb{C}^{\alpha}_0$, $0 \leqslant t < T \leqslant T_0$, and $k \in [2, \sqrt{L(b)}/L(\sigma)^2]$.

The following estimate plays a key role. **Proposition 6.** Let $u = \{u(t, x)\}_{t \ge 0, x \in [0,1]}$ denote the solution of SHE, and define $\varpi := \max(12, 6/\alpha)$ and fix $T_0 > 0$. If $u_0 \in \mathbb{C}_0^{\alpha}$ for some $\alpha \in (0, 1]$ and $\sqrt{\operatorname{Lip}(b)} > \operatorname{\varpi}\operatorname{Lip}(\sigma)^2$, then there exists a finite constant *A*—independent of L(b), $L(\sigma)$ —such that for all $T \in [0, T_0]$,

$$E\left(\sup_{t\in[0,T]}\sup_{x\in[0,1]}|u(t,x)|^{k}\right) \leqslant A^{k}(1\vee T)^{k(1+\frac{\alpha}{2}\wedge\frac{1}{4})}\left(\|u_{0}\|_{\mathbb{C}_{0}^{\alpha}}^{k}\right)$$

$$+k^{k/2}\mathcal{M}_{1}^{k}+k^{k/2}\mathcal{M}_{2}^{k}\mathcal{M}_{3}^{k}e^{AkL(b)T}\right),$$

$$(17)$$

uniformly for all $k \in \left(\varpi, \sqrt{L(b)}/L(\sigma)^2 \right]$.

The proof of Proposition 6 is also lengthy. It requires estimates providing precise dependence of the moment bounds of

$$E[|u(t,x)-u(s,y)|^k]$$

on the Linear growth constants of the coefficients b and σ .

For all $N \ge 1$ let b_N be the following truncation of the drift function:

$$b_N(z) := \begin{cases} b(z) & \text{if } |z| \leq N, \\ b(N) & \text{if } z > N, \\ b(-N) & \text{if } z < -N. \end{cases}$$
(18)

Let $\sigma_N(z)$ denote the corresponding truncation of the diffusion coefficient σ .

Consider the stochastic PDE

$$\dot{u}_{N}(t,x) = \frac{1}{2}u_{N}''(t,x) + b_{N}(u_{N}(t,x)) + \sigma_{N}(u_{N}(t,x))\xi(t,x),$$
(19)

subject to $u_N(0) = u_0$. Since b_N is globally Lipschitz, the solution u_N exists for all time.

Sketch of the proof of Theorem 2

Consider also the stopping times

$$\tau_N := \inf \left\{ t > 0: \sup_{x \in [0,1]} |u_N(t,x)| > N \right\},$$

where $\inf \varnothing := \infty$. One has

$$u_N(t,x) = u_{N+1}(t,x)$$
 for all $t \in [0, \tau_N)$ and $x \in [0,1]$.

Since u_N is well defined for all time, and is a continuous function of (t, x), this proves that $\tau_N \leq \tau_{N+1}$ a.s. for all $N \geq 1$, and therefore there exists a space-time stochastic process u such that for all $N \geq 1$, $u(t, x) = u_N(t, x)$ for all $x \in [0, 1]$ and $t \in [0, \tau_N)$. Consider the stopping time

$$\tau_{\infty} = \lim_{N \uparrow \infty} \tau_N.$$

The aim is to show that $\tau_{\infty} = \infty$ a.s.

The proof is divided into several steps. We first assume that the drift *b* in (1) has the following special form: There exist two constants $\vartheta_1, \vartheta_2 \in \mathbb{R}$ such that $\vartheta_2 \neq 0$ and

$$ilde{b}(z) = artheta_1 + artheta_2 |z| \log_+ |z| \qquad ext{for all } z \in \mathbb{R},$$
 (20)

where we recall $\log_+(a) := \log(a \lor e)$ for all $a \ge 0$.

Define

$$\widetilde{b}_{N}(z) := \vartheta_{1} + \vartheta_{2}(|z| \wedge N) \log_{+}(|z| \wedge N),$$

for all $N \ge 3$. We can take

$$L(\tilde{b}_N) = \vartheta_2(\log N). \tag{21}$$

Sketch of the proof of Theorem 2

For every fixed integer $N \ge 3$, the following stochastic PDE is well posed for all time:

$$\dot{U}_{\mathcal{N}}(t,x) = rac{1}{2}U_{\mathcal{N}}''(t,x) + \widetilde{b}_{\mathcal{N}}\left(U_{\mathcal{N}}(t,x)
ight) + \sigma_{\mathcal{N}}\left(U_{\mathcal{N}}(t,x)
ight) \xi(t,x),$$

valid for all t > 0 and $x \in [0, 1]$, subject to $U_N(0) \equiv u_0$. Define

$$au_N^{(1)} := \inf \left\{ t > 0 : \sup_{x \in [0,1]} |U_N(t,x)| > N \right\},$$

where $\inf \varnothing := \infty.$ As an important part of the proof, we need to show that

$$\tau_{\infty}^{(1)} := \lim_{N \nearrow \infty} \tau_N^{(1)} = \infty \qquad \text{a.s.}$$
(22)

We apply Proposition 6 to show that $\tau_{\infty}^{(1)}$ is greater than a positive, deterministic constant δ . In order to justify this assertion, we appeal to the Chebyshev inequality to see that for every $\varepsilon > 0$ and $N \ge 3$,

$$P\left\{\tau_{N}^{(1)} < \varepsilon\right\} = P\left\{\sup_{t \in [0,\varepsilon]} \sup_{x \in [0,1]} |U_{N}(t,x)| > N\right\}$$
$$\leqslant \quad N^{-k} E\left(\sup_{t \in [0,\varepsilon]} \sup_{x \in [0,1]} |U_{N}(t,x)|^{k}\right).$$
(23)

Next, we may apply (21) and Proposition 6 in order to see that there exist universal constants A and B such that

$$E\left(\sup_{t\in[0,\varepsilon]}\sup_{x\in[0,1]}|U_N(t,x)|^k\right) \leqslant A^k \|u_0\|_{\mathbb{C}_0^{\alpha}}^k (B+\log N)^k N^{Ak\vartheta_2\varepsilon}.$$
(24)
Here we have used the assumption that $\sigma(z) = o(|z|(\log|z|)^{\frac{1}{4}})$ as

 $|z| \rightarrow \infty$ in order to be able to apply Proposition 6.

In other words, we now have

$$P\{\tau_N^{(1)} < \varepsilon\} \leqslant A^k \|u_0\|_{\mathbb{C}_0^{\alpha}}^k (B + \log N)^k N^{k(A\vartheta_2\varepsilon - 1)}, \qquad (25)$$

uniformly for all sufficiently large integers N and $\varepsilon \in (0, 1)$. Provided that $\varepsilon < A^{-1}\vartheta_2^{-1}$, the right-hand side converges to 0 as $N \to \infty$, so (25) implies that $\tau_{\infty}^{(1)} \ge \varepsilon$ with probability one. This in turn proves that

$$au_{\infty}^{(1)} > \delta := rac{1}{2} \min(A^{-1} \vartheta_2^{-1}, 1)$$
 a.s. (26)

As δ is independent of the initial function, we can further exploit the Markov property to prove that $\tau_{\infty}^{(1)} = \infty$.

Sketch of the proof of Theorem 2

Finally we prove the theorem in the general case where *b* is an arbitrary locally-Lipschitz function that satisfies the growth condition $|b(z)| = O(|z| \log |z|)$ as $|z| \to \infty$. We can find $\vartheta_1 \in \mathbb{R}$ and $\vartheta_2 > 0$ such that

$$b_-(z)\leqslant b(z)\leqslant b_+(z),\qquad ext{for all }z\in\mathbb{R},$$

where

$$b_{\pm}(z) := \vartheta_1 \pm \vartheta_2 |z| \log_+ |z|, \quad \text{for all } z \in \mathbb{R}.$$

Let $U_{\pm}(t, x)$ denote the solution to (1), where *b* is replaced by b_{\pm} . By analogy with (18), let $b_{N,-}$ and $b_{N,+}$ be the truncations of b_{-} and b_{+} , respectively. Then

$$b_{N,-}(z) \leqslant b_N(z) \leqslant b_{N,+}(z).$$

Let u_N be the solution to (19), $U_{N,-}$ (resp. $U_{N,+}$) be the solution to (19) with b_N replaced by $b_{N,-}$ (resp. $b_{N,+}$). According to the comparison theorem, for all $(t, x) \in \mathbb{R}_+ \times [0, 1]$,

$$U_{N,-}(t,x) \leqslant u_N(t,x) \leqslant U_{N,+}(t,x). \tag{27}$$

We have shown in Step 2 that

$$\sup_{t \in [0,T]} \sup_{x \in [0,1]} |U_{\pm}(t,x)| < \infty \quad \text{ for all } T > 0.$$
 (28)

For any given (t, x), for N sufficiently large, $U_{\pm}(t, x) = U_{N,\pm}(t, x)$, therefore (27) implies that

$$U_{-}(t,x) \leqslant u_{N}(t,x) \leqslant U_{+}(t,x).$$
⁽²⁹⁾

Recall that

$$\tau_N = \inf\{t > 0 : \sup_{x \in [0,1]} |u_N(t,x)| > N\}.$$

Then (28) and (29) imply that $\lim_{N\to\infty}\tau_N=\infty$ a.s., and we can define

$$u(t,x) = u_N(t,x),$$
 for $t \in [0,\tau_N]$ and $x \in [0,1].$

As above, this definition is coherent. By (29),

$$U_{-}(t,x)\leqslant u(t,x)\leqslant U_{+}(t,x), \qquad ext{for all }t\in \mathbb{R}_{+} ext{ and }x\in [0,1].$$

We can show that u is the solution of the equation:

$$\begin{split} u(t,x) &= (\mathcal{G}_t u_0)(x) + \int_{(0,t)\times(0,1)} \mathcal{G}_{t-s}(x,y) b(u(s,y)) \, \mathrm{d}s \, \mathrm{d}y \\ &+ \int_{(0,t)\times(0,1)} \mathcal{G}_{t-s}(x,y) \sigma(u(s,y)) \, \xi(\mathrm{d}s \, \mathrm{d}y), \end{split}$$

and

$$\sup_{t\in[0,T]}\sup_{x\in[0,1]}|u(t,x)|<\infty\qquad\text{for all }T>0,$$

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