

Atomistic/continuum coupling schemes for solids
(a priori error analysis of a prototypical energy-based coupling scheme)

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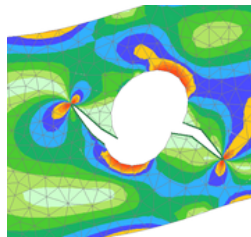
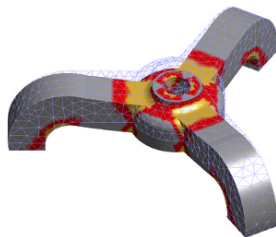
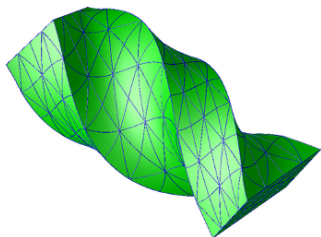
Continuum Solid Mechanics (Elasticity)

- Reference domain $\Omega \subset \mathbb{R}^d$, Deformation $y : \Omega \rightarrow \mathbb{R}^d$
- Minimize **stored energy**

$$\min \mathcal{E}^c(y) := \int_{\Omega} [W(Dy) - g \cdot y] dx, \quad y \in \mathcal{Y}.$$

- Galerkin discretization: choose $\mathcal{Y}_h \subset \mathcal{Y}$, $\dim \mathcal{Y}_h < +\infty$,

$$\min \mathcal{E}^c(y_h), \quad y_h \in \mathcal{Y}_h.$$



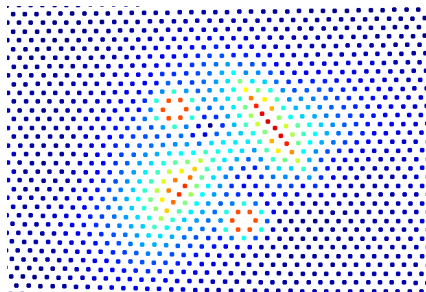
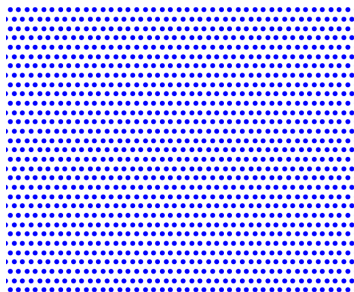
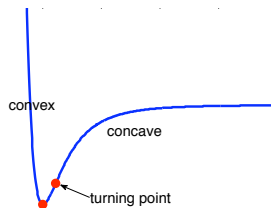
images taken from GETFEM++ website

Atomistic Mechanics

- **Atomistic body** with N atoms at positions $\mathbf{y} = (y_n)_{n=1}^N \in \mathbb{R}^{d \times N}$
- **Total energy** of configuration \mathbf{y} :

$$\mathcal{E}^a(\mathbf{y}) = \sum_{n \neq k} \phi(|y_n - y_k|) - \langle \mathbf{g}, \mathbf{y} \rangle$$

- ϕ : Lennard–Jones type potential
- \mathbf{g} : external forces;
for simplicity assume dead loads
- **Find local minimizer of $\mathcal{E}^a(\mathbf{y})$ over all admissible configurations \mathbf{y} !**



Overview

1 The simplest A/C hybrid energy:

- 1D model problem
- energy-based QC method (*brutal energy mixing*)
- the problem of ghost forces
- ghost forces vs. consistency, part 1

2 Energy-based ghost-force removal; Consistency:

- quasinonlocal QC (*bond-based coupling*)
- ghost force vs. consistency, part 2 & 3
- outlook to 2D/3D

3 Stability & Error estimates:

- abstract framework
- stability analysis via strain gradients
- a rigorous error estimates

A 1D Model Problem

Periodic displacements:

$$\mathcal{U} = \{ \mathbf{u} = (u_n)_{n \in \mathbb{Z}} : u_{n+N} = u_n, \sum_{n=1}^N u_n = 0 \},$$
$$\mathcal{Y} = \{ \mathbf{y} = (y_n)_{n \in \mathbb{Z}} : y_n = x_n + u_n \text{ where } \mathbf{u} \in \mathcal{U} \}.$$

Atomistic energy: (*next-nearest neighbour pair interactions*)

$$\mathcal{E}^a(\mathbf{y}) = \varepsilon \sum_{n=1}^N \phi(y'_n) + \varepsilon \sum_{n=1}^N \phi(y'_n + y'_{n+1}) = \varepsilon \sum_{n=1}^N \mathcal{E}_n^a(\mathbf{y}),$$

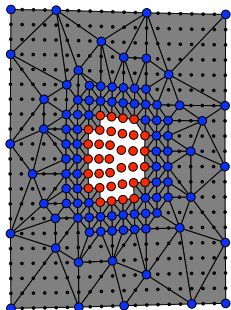
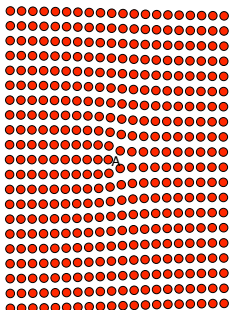
where $\mathcal{E}_n^a(\mathbf{y}) = \frac{1}{2} \{ \phi(y'_{n-1} + y'_n) + \phi(y'_n) + \phi(y'_{n+1}) + \phi(y'_n + y'_{n+1}) \}.$

Continuum finite element model with $h = \varepsilon$:

$$\mathcal{E}^c(\mathbf{y}) = \varepsilon \sum_{n=1}^N \{ \phi(y'_n) + \phi(2y'_n) \} = \varepsilon \sum_{n=1}^N \mathcal{E}_n^c(\mathbf{y})$$

where $\mathcal{E}_n^c(\mathbf{y}) = \frac{1}{2} \{ \phi(2y'_n) + \phi(y'_n) + \phi(y'_{n+1}) + \phi(2y'_n) \}.$

The Energy-Based Quasicontinuum Method



- Choose atomistic and continuum regions:

$$\mathcal{N}^a \cup \mathcal{N}^c = \{1, \dots, N\}$$

- Define a/c hybrid energy (“brutal energy mixing”)

$$\mathcal{E}^{\text{qc}}(\mathbf{y}) = \varepsilon \sum_{n \in \mathcal{N}^a} \mathcal{E}_n^a(\mathbf{y}) + \varepsilon \underbrace{\sum_{n \in \mathcal{N}^c} \mathcal{E}_n^c(\mathbf{y})}_{\int_{\Omega^c} W(D\mathbf{y}) \, dx} - \langle \mathbf{g}, \mathbf{y} \rangle$$



[Tadmor, Miller, Ortiz; 1996]

The Ghost Forces = Patch Test (1)

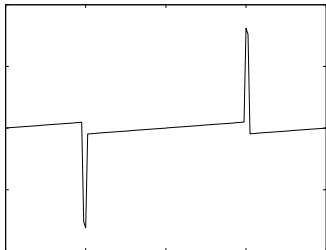
- $\mathbf{g} = 0 \Rightarrow$ exact solution is the reference lattice: $\mathbf{y}^a = \mathbf{x}$.
- Continuum model is exact here: $\nabla \mathcal{E}^a(\mathbf{x}) = \nabla \mathcal{E}^c(\mathbf{x}) = 0$
- Patch test for $\mathcal{E}^{\text{qc}}(\mathbf{y}) = \varepsilon \sum_{\mathcal{N}^a} \mathcal{E}_n^a + \varepsilon \sum_{\mathcal{N}^c} \mathcal{E}_n^c$:

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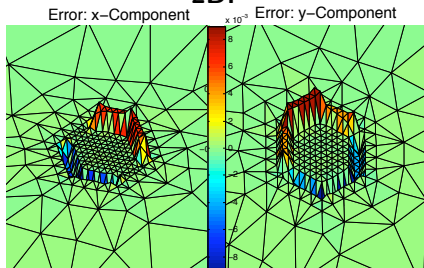
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- Patch test for $\mathcal{E}^{\text{qc}}(\mathbf{y}) = \varepsilon \sum_{\mathcal{N}^a} \mathcal{E}_n^a + \varepsilon \sum_{\mathcal{N}^c} \mathcal{E}_n^c$:

The ghost forces: $\nabla \mathcal{E}^{\text{qc}}(\mathbf{x}) \neq 0$ at the interface

1D:



2D:



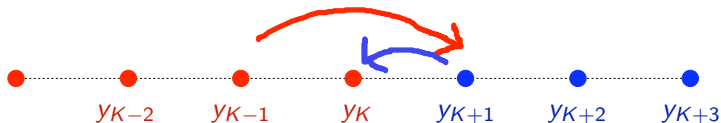
The Ghost Forces = Patch Test (2)

Solutions for \mathcal{E}^a and \mathcal{E}^c :

$$\nabla \mathcal{E}^a(\mathbf{x}) = 0 \quad \text{and} \quad \nabla \mathcal{E}^c(\mathbf{x}) = 0$$

But if we insert $\mathbf{y}_a = \mathbf{x}$ into $\nabla \mathcal{E}^{\text{qc}}$: (“consistency error”)

$$\frac{\partial \mathcal{E}^{\text{qc}}}{\partial y_n} \Big|_{\mathbf{y}=\mathbf{x}} = \frac{\phi'(2)}{2} \times \begin{cases} 0, & n = \dots, K-2 \\ 1, & n = K-1 \\ -1, & n = K \\ -1, & n = K+1 \\ 1, & n = K+2 \\ 0, & n = K+3, \dots \end{cases}$$



Ghost Force vs. Consistency, Part 1

Theorem

Ghost forces \Rightarrow high consistency error \Rightarrow low accuracy.

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Proof.

1. Observation: $\nabla_{L^2} = \varepsilon^{-1} \nabla_{\ell^2} = \varepsilon^{-1} \nabla =: \nabla_\varepsilon$
2. Consistency Error = **a2c Model Error** + Coupling Error
+ ~~Discretization Error~~
3. Insert calculation from previous page:

$$\|\nabla_\varepsilon \mathcal{E}^{\text{qc}}(\mathbf{y}^{\text{a}})\|_{L^\infty} \sim \varepsilon^{-1}$$

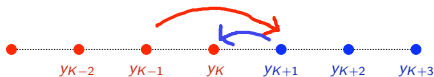
$$\Rightarrow \|\nabla_\varepsilon \mathcal{E}^{\text{qc}}(\mathbf{y}^{\text{a}})\|_{W^{-1,\infty}} \gtrsim 1$$

$$\Rightarrow \|\mathbf{y}^{\text{a}} - \mathbf{y}^{\text{qc}}\|_{W^{1,\infty}} \gtrsim 1.$$

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Ghost Force Removal



Observation: ghost forces can arise from asymmetry in the interaction

Idea: concentrate on bonds rather than atoms.

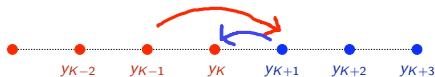
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$$\mathcal{E}^c(\mathbf{y}) = \varepsilon \sum_n \phi(y'_n) + \varepsilon \sum_n \phi(2y'_n) = \int W(D\mathbf{y}) \, dx$$

Here, going from atomistic to continuum can be understood as follows:

$$\phi(y'_n + y'_{n+1}) \approx \frac{1}{2} \{ \phi(2y'_n) + \phi(2y'_{n+1}) \}$$

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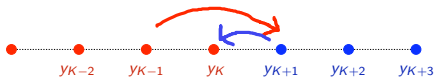
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↪ **new A/C hybrid energy:**

$$\mathcal{E}^{qc}(\mathbf{y}) = \varepsilon \sum_n \phi(y'_n) + \varepsilon \sum_{n \in \mathcal{N}^a} \phi(y'_n + y'_{n+1}) + \varepsilon \sum_{n \in \mathcal{N}^c} \left\{ \frac{1}{2} \phi(2y'_n) + \frac{1}{2} \phi(2y'_{n+1}) \right\}$$

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$$= \varepsilon \sum_{n \in \mathcal{N}^a} \mathcal{E}_n^a + \int_{\Omega^c} W(D\mathbf{y}) dx + \text{interface correction}$$

Ghost Forces vs. Consistency, Part 2

Easy to check that $\nabla \mathcal{E}^{\text{qc}}(\mathbf{x}) = 0$; does this imply higher accuracy?

Theorem

$$\|\nabla_{\varepsilon} \mathcal{E}^{\text{a}}(\mathbf{y}) - \nabla_{\varepsilon} \mathcal{E}^{\text{qc}}(\mathbf{y})\|_{W^{-1,p}} \lesssim \varepsilon \|\mathbf{y}''\|_{L^p(\mathcal{N}^i)} + \varepsilon^2 \|\mathbf{y}'''\|_{L^p(\mathcal{N}^c)} + \varepsilon^2 \|\mathbf{y}''\|_{L^{2p}(\mathcal{N}^c)}^2$$

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Proof. Fix $\mathbf{u} \in \mathcal{U}$, bound $\langle \nabla_{\varepsilon} \mathcal{E}^{\text{a}} - \nabla_{\varepsilon} \mathcal{E}^{\text{qc}}, \mathbf{u} \rangle_{\varepsilon}$ above.

1. Calculate Fréchet derivative

$$\begin{aligned} \langle \nabla_{\varepsilon} \mathcal{E}^{\text{a}} - \nabla_{\varepsilon} \mathcal{E}^{\text{qc}}, \mathbf{u} \rangle_{\varepsilon} &= \varepsilon \sum_{n \in \mathcal{N}^c} \{ \phi'(y'_n + y'_{n+1})[u'_n + u'_{n+1}] \\ &\quad - \phi'(2y'_n)[u'_n] - \phi'(2y'_{n+1})[u'_{n+1}] \} \\ &= \varepsilon \sum_{n \in \mathcal{N}^c} T_n u'_n \leq \|T_n\|_{L^p} \|\mathbf{u}'\|_{L^{p'}} \end{aligned}$$

2. Estimate T_n using Taylor expansion:

$$T_n = \begin{cases} \phi'(y'_{n-1} + y'_n) + \phi'(y'_n + y'_{n+1}) - 2\phi'(2y'_n), & n \in \text{interior}(\mathcal{N}^c) \\ \phi'(y'_n + y'_{n+1}) - \phi'(2y'_n), & n \in \text{a/c interface.} \end{cases}$$

Ghost Forces vs. Consistency, Part 3

- Remarks:**
1. In 1D these results are generic! In 2D/3D it is unclear whether absence of ghost forces gives improved accuracy.
 2. $\varepsilon = h$ is not necessary; analysis can admit general meshes.

Ghost Forces vs. Consistency, Part 3

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1. In 1D these results are generic! In 2D/3D it is unclear whether absence of ghost forces gives improved accuracy.
 2. $\varepsilon = h$ is not necessary; analysis can admit general meshes.

■ Local QC / Cauchy–Born:

$$\|\nabla_\varepsilon \mathcal{E}^a - \nabla_\varepsilon \mathcal{E}^c\|_{W_h^{-1,p}} \lesssim \|h\mathbf{y}_a''\|_{L^p}^2 + \dots$$

■ QCE with ghost force:

$$\|\nabla_\varepsilon \mathcal{E}^a - \nabla_\varepsilon \mathcal{E}^{\text{qc}}\|_{W_h^{-1,p}} \lesssim \|h\mathbf{y}_a''\|_{L^p(\mathcal{N}^c)}^2 + \varepsilon \|\mathbf{y}_a''\|_{L^p(\mathcal{N}^i)} + \varepsilon^{1/p} + \dots$$

■ QCE without ghost force:

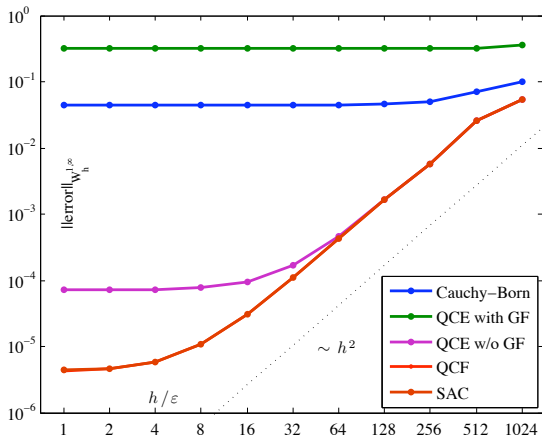
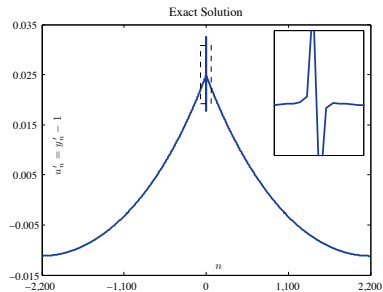
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■ Force-based QC: *(not discussed in this talk)*

$$\|\nabla_\varepsilon \mathcal{E}^a - F^{\text{qc}}\|_{W_h^{-1,\infty}} \lesssim \|h\mathbf{y}_a''\|_{L^\infty(\mathcal{N}^c)}^2 + \dots$$

1D Numerical Experiment

$$N = 1024, \phi(r) = \exp(-2\alpha(r - 1)) - 2 \exp(-\alpha(r - 1))$$



Extension to 2D

- **Crucial idea in 1D:**

$$\phi(y'_n + y'_{n+1}) \approx \frac{1}{2} \{ \phi(2y'_n) + \phi(2y'_{n+1}) \} = \frac{1}{2} \int_{x_{n-1}}^{x_{n+1}} \phi(y'(x)) dx$$

- **Natural generalization to dD:** identify \mathbf{y} with P1-interpolant $\bar{y}(x)$ then

$$\phi(|y(\xi + \rho) - y(\xi)|) \approx \int_{t=0}^1 \phi(|Dy(\xi + t\rho)\rho|) dt$$

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- **Construction of QC method:**

[Shapeev; preprint]

$$\begin{aligned} & \sum_{(\xi, \xi+\rho)} \phi(|\bar{y}(\xi + \rho) - \bar{y}(\xi)|) \\ \approx & \sum_{(\xi, \xi+\rho) \in \mathcal{B}_a} \phi(|\bar{y}(\xi + \rho) - \bar{y}(\xi)|) + \sum_{(\xi, \xi+\rho) \in \mathcal{B}_c} \int_{t=0}^1 \phi(|D\bar{y}(\xi + t\rho)\rho|) dt \\ \stackrel{2D}{=} & \sum_{(\xi, \xi+\rho) \in \mathcal{B}_a} \phi(|\bar{y}(\xi + \rho) - \bar{y}(\xi)|) + \int_{\Omega_c} W(D\bar{y}) dx + \text{interface correction} \end{aligned}$$

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- **Consistency error estimate:**

[O. & Shapeev; in prep.]

$$\langle \nabla_\varepsilon \mathcal{E}^a(\mathbf{y}) - \nabla_\varepsilon \mathcal{E}^{qc}(\mathbf{y}), \mathbf{u} \rangle \lesssim \left\{ \inf_{\substack{\tilde{y} \in C^2 \\ \tilde{y}(\xi) = y_\xi}} \|hD^2\tilde{y}\|_{\tilde{L}^p(\Omega_c)} \right\} \|D\bar{u}\|_{L^{p'}(\Omega_c)}$$

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Abstract Framework

Assumptions on an exact solution:

- *Stability*: \mathbf{y}_a is a stable equilibrium of \mathcal{E}^a :

$$\nabla_{\varepsilon} \mathcal{E}^a(\mathbf{y}_a) = 0 \quad \text{and} \quad \nabla_{\varepsilon}^2 \mathcal{E}^a(\mathbf{y}_a) > 0$$

- *Smoothness*: \mathbf{y}_a is “smooth” in \mathcal{N}^c .

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Outline of an Error Analysis:

1. *Consistency*: $\nabla_{\varepsilon} \mathcal{E}^{\text{qc}}(\mathbf{y}_a) = \nabla_{\varepsilon} \mathcal{E}^{\text{qc}}(\mathbf{y}_a) - \nabla_{\varepsilon} \mathcal{E}^a(\mathbf{y}_a)$ is small
2. *Stability*: can we prove $\nabla_{\varepsilon}^2 \mathcal{E}^{\text{qc}}(\mathbf{y}_a) > 0$??
3. Inverse function theorem \Rightarrow

$$\exists \mathbf{y}_{\text{qc}} \text{ s.t. } \nabla_{\varepsilon} \mathcal{E}^{\text{qc}}(\mathbf{y}_{\text{qc}}) = 0, \quad \text{and}$$

$$\|\mathbf{y}'_{\text{qc}} - \mathbf{y}'_a\|_{L^p} \lesssim \|\nabla_{\varepsilon} \mathcal{E}^{\text{qc}}(\mathbf{y}_a) - \nabla_{\varepsilon} \mathcal{E}^a(\mathbf{y}_a)\|_{W^{-1,p}}$$

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Task: Understand relationship between

$$\alpha^a(\mathbf{y}) = \inf_{\|\mathbf{u}'\|_{L^2}=1} \langle \nabla_{\varepsilon}^2 \mathcal{E}^a(\mathbf{y}) \mathbf{u}, \mathbf{u} \rangle \quad \text{and} \quad \alpha^{\text{qc}}(\mathbf{y}) = \inf_{\|\mathbf{u}'\|_{L^2}=1} \langle \nabla_{\varepsilon}^2 \mathcal{E}^{\text{qc}}(\mathbf{y}) \mathbf{u}, \mathbf{u} \rangle$$

Hessians & Strain Gradients (1)

Hessian of continuum energy:

$$\langle \nabla_{\varepsilon}^2 \mathcal{E}^c(\mathbf{y}) \mathbf{u}, \mathbf{u} \rangle = \varepsilon \sum_n \phi''(y'_n) |u'_n|^2 + \varepsilon \sum_n 4\phi''(2y'_n) |u'_n|^2$$

Hessian of atomistic energy:

$$\langle \nabla_{\varepsilon}^2 \mathcal{E}^a(\mathbf{y}) \mathbf{u}, \mathbf{u} \rangle = \varepsilon \sum_n \phi''(y'_n) |u'_n|^2 + \varepsilon \sum_{n=1}^N \phi''(y'_n + y'_{n+1}) |u'_n + u'_{n+1}|^2$$

Hessian of QC energy:

$$\begin{aligned} \langle \nabla_{\varepsilon}^2 \mathcal{E}^{qc}(\mathbf{y}) \mathbf{u}, \mathbf{u} \rangle &= \varepsilon \sum_n \phi''(y'_n) |u'_n|^2 + \varepsilon \sum_{n \in \mathcal{N}^a} \phi''(y'_n + y'_{n+1}) |u'_n + u'_{n+1}|^2 \\ &\quad + \varepsilon \sum_{n \in \mathcal{N}^c} \{2\phi''(2y'_n) |u'_n|^2 + 2\phi''(2y'_{n+1}) |u'_{n+1}|^2\} \end{aligned}$$

Hessians & Strain Gradients (1)

Hessian of continuum energy:

$$\langle \nabla_{\varepsilon}^2 \mathcal{E}^c(\mathbf{y}) \mathbf{u}, \mathbf{u} \rangle = \varepsilon \sum_n \phi''(y'_n) |u'_n|^2 + \varepsilon \sum_n 4\phi''(2y'_n) |u'_n|^2$$

Hessian of atomistic energy:

$$\langle \nabla_{\varepsilon}^2 \mathcal{E}^a(\mathbf{y}) \mathbf{u}, \mathbf{u} \rangle = \varepsilon \sum_n \phi''(y'_n) |u'_n|^2 + \varepsilon \sum_{n=1}^N \phi''(y'_n + y'_{n+1}) |u'_n + u'_{n+1}|^2$$

Hessian of QC energy:

$$\begin{aligned} \langle \nabla_{\varepsilon}^2 \mathcal{E}^{qc}(\mathbf{y}) \mathbf{u}, \mathbf{u} \rangle &= \varepsilon \sum_n \phi''(y'_n) |u'_n|^2 + \varepsilon \sum_{n \in \mathcal{N}^a} \phi''(y'_n + y'_{n+1}) |u'_n + u'_{n+1}|^2 \\ &\quad + \varepsilon \sum_{n \in \mathcal{N}^c} \{2\phi''(2y'_n) |u'_n|^2 + 2\phi''(2y'_{n+1}) |u'_{n+1}|^2\} \end{aligned}$$

Rewrite NNN terms using the *Polarisation Identity*:

$$|u'_n + u'_{n+1}|^2 = 2|u'_n|^2 + 2|u'_{n+1}|^2 - \varepsilon^2 |u''_n|^2$$

Hessians & Strain Gradients (2)

$$\langle \nabla_{\varepsilon}^2 \mathcal{E}^c(\mathbf{y}) \mathbf{u}, \mathbf{u} \rangle = \varepsilon \sum_n A_n^c |u'_n|^2$$

$$\langle \nabla_{\varepsilon}^2 \mathcal{E}^a(\mathbf{y}) \mathbf{u}, \mathbf{u} \rangle = \varepsilon \sum_n A_n^a |u'_n|^2 + \varepsilon^3 \sum_n B_n |u''_n|^2$$

$$\langle \nabla_{\varepsilon}^2 \mathcal{E}^{\text{qc}}(\mathbf{y}) \mathbf{u}, \mathbf{u} \rangle = \varepsilon \sum_n A_n^{\text{qc}} |u'_n|^2 + \varepsilon^3 \sum_{n \in \mathcal{N}^a} B_n |u''_n|^2$$

where $A_n^a \approx A_n^{\text{qc}}$ and $B_n \geq 0$

Theorem

- 1 $\alpha^a(\mathbf{y}) \geq \alpha^{\text{qc}}(\mathbf{y}) - C\varepsilon \|\mathbf{y}''\|_{L^\infty(\mathcal{N}^c)}$
- 2 $\alpha^{\text{qc}}(\mathbf{y}) \geq \min_n A_n^{\text{qc}}$
- 3 If \mathbf{y} is globally smooth then $\alpha^{\text{qc}} \geq \alpha^a - O(\varepsilon)$

Error Estimate

Theorem (Existence and A Priori Error Estimate)

Let \mathbf{y}_a be a stable equilibrium of \mathcal{E}^a s.t. $\min_n A_n^{\text{qc}} > 0$. Then there exists a constant $\delta > 0$ such that, if

$$\varepsilon \|\mathbf{y}''\|_{L^\infty(\mathcal{N}^c)} + \varepsilon^{3/2} \|\mathbf{y}'''\|_{L^2(\mathcal{N}^c)} + \varepsilon^{3/2} \|\mathbf{y}''\|_{L^4(\mathcal{N}^c)}^2 \leq \delta,$$

then there exists a stable equilibrium \mathbf{y}_{qc} of \mathcal{E}^{qc} s.t.

$$\|\mathbf{y}'_{\text{qc}} - \mathbf{y}'_a\|_{L^2} \lesssim \frac{1}{\alpha^{\text{qc}}(\mathbf{y}_a)} \left[\varepsilon \|\mathbf{y}''_a\|_{L^2(\mathcal{N}^i)} + \varepsilon^2 (\|\mathbf{y}'''\|_{L^2(\mathcal{N}^c)} + \|\mathbf{y}''_a\|_{L^4(\mathcal{N}^c)}^2) \right]$$

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$$\|\mathbf{y}'_{\text{qc}} - \mathbf{y}'_a\|_{L^2} \lesssim \frac{1}{\alpha^{\text{qc}}(\mathbf{y}_a)} \left[\varepsilon \|\mathbf{y}''_a\|_{L^2(\mathcal{N}^i)} + \varepsilon^2 (\|\mathbf{y}'''\|_{L^2(\mathcal{N}^c)} + \|\mathbf{y}''_a\|_{L^4(\mathcal{N}^c)}^2) \right]$$

Remarks:

- Can admit defects in the analysis, but need to modify the function space setting. [O. & Süli; 2008], [O.; 2010]
- Sharp stability in 2D/3D version is difficult (work in progress).

Atomistic/Continuum Hybrid Methods: POV

1 Direct Coarse Graining without Coupling

- Force-based: [Knap & Ortiz; 2003], Energy-based: [Eidel & Stuchowsky; 2008]
- Analysis: [Lin; 2003, 2008], [O. & Süli; 2008], [Luskin & O.; 2009]
- Conclusion so far: not as simple as it sounds ...

2 Energy Based Coupling

- Original QC (with GF): [Tadmor & Ortiz & Phillips; 1996]
- Ghost force removal: [Shimokawa et al; 2004], [E & Lu & Yang; 2006], [Shapeev]
- Analysis: 1D now fairly advanced; 2D in progress
- Still some unresolved difficulties general interaction potentials ...

3 Force Based Coupling

- Naturally free of ghost forces
- Large number of different variants: Schmauder, Gumbsch; Shenoy, Rodney, Ortiz, Tadmor; Miller, Curtin; Luskin, Dobson; Parks, Bochev, Lehoucq; ...
- Beautiful (and difficult) question open in their analysis.

Other Methods: Blending, Domain Decomposition, various methods for finite temperature, dynamics, DFT: Parks, Gunzburger, Seleson, Lehoucq, Karpov, Park, Rudd, Broughton, Ortiz, Gavini, Bhattacharya, E, Lu, Garcia-Cervera, Cancès, Deleurence, Lewin, Legoll, ...