An Eulerian finite element method for elliptic equations on moving surfaces



levitated drop fluid/fluid



transport of surfactants

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Motivation: Two-phase fluid dynamics + surfactants

- Standard model for two-phase flow.
- Level set for interface capturing.
- Treatment of surface tension.
- Special FE space for pressure.

An Eulerian FEM for elliptic equations on moving surfaces

• FEM for transport equation on the interface







$$\begin{cases} \rho_i (\mathbf{u}_t + (\mathbf{u} \cdot \nabla) \mathbf{u}) = \operatorname{div} (\sigma) + \rho_i \mathbf{g} \\ = -\nabla p + \operatorname{div} (\mu_i \mathbf{D}(\mathbf{u})) + \rho_i \mathbf{g} & \text{in } \Omega_i & \text{for } i = 1, 2 \\ \operatorname{div} \mathbf{u} = 0 & \text{in } \Omega_i \\ [\sigma \mathbf{n}]_{\Gamma} = \tau \mathcal{K} \mathbf{n} - \nabla_{\Gamma} \tau, \quad [\mathbf{u}]_{\Gamma} = 0 . \end{cases}$$

 $D(\mathbf{u}) = \nabla \mathbf{u} + \nabla \mathbf{u}^T$:deformation tensor \mathcal{K} :curvature of Γ $\sigma = -p \mathbf{I} + \mu \mathbf{D}(\mathbf{u})$:stress tensorAssumption: τ constantAppropriate model for a large class of two-phase flow problems.





Idea:(Sethian, Osher) $\Gamma(t) =$ zero-level of a scalar function The level set function $\varphi(x,t)$

$$\varphi(x,t) = \begin{cases} < 0 & \text{for } x \text{ in phase } \Omega_1 \\ > 0 & \text{for } x \text{ in phase } \Omega_2 \\ = 0 & \text{at the interface} \end{cases}$$



should be an *"approximate signed distance function"*.

$$x(t) \in \Gamma(t) \Rightarrow \varphi(x(t), t) = 0.$$

Level set equation

$$\varphi_t + \mathbf{u} \cdot \nabla \varphi = 0$$





Navier-Stokes equations coupled with level set equation:

$$\rho(\varphi) \Big(\mathbf{u}_t + (\mathbf{u} \cdot \nabla) \mathbf{u} \Big) - \operatorname{div} \Big(\mu(\varphi) \, \mathbf{D}(\mathbf{u}) \Big) + \nabla p = \rho(\varphi) \, g - \tau \, \mathcal{K}(\varphi) \, \delta_{\Gamma} \mathbf{n}_{\Gamma}$$
$$\nabla \cdot \mathbf{u} = 0$$
$$\varphi_t + \mathbf{u} \cdot \nabla \varphi = 0$$

where ρ, μ and $\mathcal{K}, \delta_{\Gamma}, \mathbf{n}_{\Gamma}$ depend on φ , e.g.:

$$\mathcal{K}(\varphi) = \nabla \cdot \left(\frac{\nabla \varphi}{||\nabla \varphi||} \right)$$
 second derivatives.

Localized force term in weak formulation:

$$f_{\Gamma}(\mathbf{v}) = \tau \int_{\Gamma} \mathcal{K} \mathbf{n}_{\Gamma} \cdot \mathbf{v} \, ds$$

$$f_{\Gamma} \in H^{-1}(\Omega)$$





Discretization:

- Weak formulation + FE methods; velocity space: P_2 .
- Discretization of localized force term f_{Γ} : Laplace-Beltrami.
- Finite element space for discontinuous pressure: XFEM $P_1 \rightsquigarrow Q_h^{\Gamma}$.
- Level set equation: P_2 FE + SDFEM stabilization.
- Time integration: θ -schema / fractional step.

Iterative solvers:

- Multigrid. Krylov subspace methods.
- Preconditioners: robustness w.r.t. μ , ρ , Δt , h.
- XFEM \longrightarrow modified iterative solvers?





 Γ = zero level of ϕ (= level set function = signed distance function) ϕ_h = piecewise quadratic FE approximation of ϕ .

Our strategy:

 $\phi \approx \phi_h$ (piecewise P_2) $\rightarrow I(\phi_h)$ (piecewise P_1 on refined mesh).

 $\Gamma \approx \Gamma_h :=$ zero level of $I(\phi_h)$ (planar segments).



Under reasonable assumptions: dist(Γ , Γ_h) $\leq c h^2$.





Belytschko (1999) for elasticity problems. Hansbo (2002) for interface problems.

Idea: Enrich FE space (e.g. P_1 FE) by additional discontinuous basis functions near Γ :

$$p_j^{\Gamma}(\mathbf{x}) := p_j(\mathbf{x}) H_{\Gamma}(\mathbf{x})$$

where
$$H_{\Gamma}(\mathbf{x}) = \begin{cases} 1 & x \in \Omega_2, \\ 0 & \text{else.} \end{cases}$$

(Technical) difficulties:

• Integration over sub-elements $T \cap \Omega_2$:

$$\int_{T} p_{j}^{\Gamma}(\mathbf{x}) \cdot f(\mathbf{x}) \, d\mathbf{x} = \int_{T \cap \Omega_{2}} p_{j}(\mathbf{x}) \cdot f(\mathbf{x}) \, d\mathbf{x}$$

- Q_h^{Γ} depends on $\Gamma!$ (in practice: Γ_h)
- Reference: [Groß, R., JCP 07].











 $f_{\Gamma}(\mathbf{v}) = \tau \int_{\Gamma} \mathcal{K} \mathbf{n}_{\Gamma} \cdot \mathbf{v} \, ds$ with $\tau = 1$. Note $\mathcal{K} = 2/r = 3$.

Solution:
$$u^* = 0$$
, $p^* = \begin{cases} C & \text{in } \Omega_2, \\ C + \tau \mathcal{K} & \text{in } \Omega_1. \end{cases}$







 Q_h^1 FE (standard P_1) Q_h^{Γ} FE (XFEM space)





ref.	$p_h \in Q_h^1$		$p_h \in Q_h^{{\sf \Gamma}_h}$	
	$\ e_p\ _{L^2}$	order	$\ e_p\ _{L^2}$	order
0	1.60E+00	_	1.64E-01	_
1	1.07E+00	0.57	4.97E-02	1.73
2	8.23E-01	0.38	1.66E-02	1.58
3	5.80E-01	0.51	7.16E-03	1.22
4	4.13E-01	0.49	2.83E-03	1.34

Pressure errors for the $P_2 - Q_h^1$ and $P_2 - Q_h^{\Gamma}$ pair.





An Eulerian FEM for elliptic equations on moving surfaces





Convection-diffusion equation for c(x,t), $x \in \Gamma(t)$:

$$\partial_{t,n}c - D_{\Gamma}\Delta_{\Gamma}c + \nabla_{\Gamma}\cdot(c\mathbf{u}_{\Gamma}) - \mathcal{K}u_{\perp}c = 0,$$

How can we discretize this equation on Γ ?.

Obvious idea: use a FE space induced by the "outer" triangulation T_h .



Define

 $\omega_h := \cup_{T \in \mathcal{F}_h} S_T : \text{ tetrahedra in } \mathcal{T}'_h \text{ intersected by } \Gamma_h$ $V_h := \{ v_h \in C(\omega_h) \mid v_{|S_T} \in P_1 \text{ for all } T \in \mathcal{F}_h \} : \text{ outer space}$

$$V_h^{\mathsf{\Gamma}} := \{\psi_h \in H^1(\mathsf{\Gamma}_h) \mid \exists v_h \in V_h : \psi_h = v_h|_{\mathsf{\Gamma}_h}\}: \text{ interface space}$$

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Laplace-Beltrami equation

$$-\Delta_{\Gamma} u + u = f \quad \text{on } \Gamma,$$

with $\Gamma=\{\mathbf{x}\in\mathbb{R}^3\mid\|\mathbf{x}\|_2=1\}$ and $\Omega=(-2,\,2)^3.$

Solution:

$$u(\mathbf{x}) = a \frac{\|\mathbf{x}\|^2}{12 + \|\mathbf{x}\|^2} \left(3x_1^2 x_2 - x_2^3 \right).$$

Tetrahedral triangulations: $\{T_l\}_{l\geq 0}$ constructed by local refinement close to Γ . Mesh size $h_l \sim \sqrt{3} 2^{-l}$.

Level set function $\phi(\mathbf{x}) = \|\mathbf{x}\|^2 - 1$; $\phi_h := I(\phi)$ piecewise linear on \mathcal{T}_h .

 $\Gamma_h := \{ \mathbf{x} \in \Omega \mid I(\phi_h)(\mathbf{x}) = \mathbf{0} \}$

Note: the interface triangulation Γ_h is not shape-regular:



Solution u:





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Determine $u_h \in V_h^{\Gamma}$ such that

$$\int_{\Gamma_h} \nabla_{\Gamma_h} u_h \nabla_{\Gamma_h} \psi_h + u_h \psi_h d\mathbf{s}_h = \int_{\Gamma_h} f_h \psi_h d\mathbf{s}_h \quad \text{for all } \psi_h \in V_h^{\Gamma},$$

with f_h an extension of f.

Results:

level l	$\ u-u_h\ _{L^2({\sf \Gamma}_h)}$	factor	
1	0.1124	—	
2	0.03244	3.47	
3	0.008843	3.67	
4	0.002186	4.05	
5	0.0005483	3.99	
6	0.0001365	4.02	
7	0.0000341	4.00	





Error analysis [Olshanskii, AR., submitted]:

Theorem. For each $u \in H^2(\Gamma)$ the following holds

$$\inf_{v_h \in V_h^{\Gamma}} \|u^e - v_h\|_{L^2(\Gamma_h)} \le \|u^e - (I_h u^e)_{|\Gamma_h}\|_{L^2(\Gamma_h)} \le C h^2 \|u\|_{H^2(\Gamma)},$$
$$\inf_{v_h \in V_h^{\Gamma}} \|u^e - v_h\|_{H^1(\Gamma_h)} \le \|u^e - (I_h u^e)_{|\Gamma_h}\|_{H^1(\Gamma_h)} \le C h \|u\|_{H^2(\Gamma)}.$$

Implementation very easy:

$$\int_{\Gamma_h} \nabla_{\Gamma_h} \phi_i \cdot \nabla_{\Gamma_h} \phi_j + \phi_i \phi_j \mathrm{ds}_h = \sum_{T \in \mathcal{F}_h} \int_T \dots \mathrm{ds}_h$$

Furthermore:

- No data structure for triangulation of Γ_h needed.
- No shape regularity of triangulation of Γ_h required.



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Let $(\phi_i)_{1 \le i \le m}$ be all nodal basis functions in V_h with support intersected by Γ_h . Then

$$\operatorname{span}\{(\phi_i)_{|\Gamma_h} \mid 1 \leq i \leq m\} = V_h^{\Gamma},$$

but $(\phi_i)_{|\Gamma_h}$ are not necessarily independent. Mass matrix:

$$M_{i,j} = \int_{\Gamma_h} \phi_i \phi_j \,\mathrm{d}\mathbf{s}_h, \quad \mathbf{1} \le i, j \le m, \quad \tilde{M} := D_M^{-\frac{1}{2}} M D_M^{-\frac{1}{2}}.$$

Spectrum of \tilde{M} :

level l	m	λ_1	λ_2	λ_m	λ_m/λ_2
1	112	3.8 e-17	0.0261	2.86	109
2	472	4.0 e-17	0.0058	2.83	488
3	1922	1.0 e-17	0.0012	2.83	2358
4	7646	3.6 e-17	0.00029	2.83	9660

Remarks:

- Scaling with D_M is essential.
- Analysis: in progress.
- Similar results for stiffness matrix.





More information:

www.igpm.rwth-aachen.de/DROPS/