Intersection of random walks in supercritical dimensions

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joint work with Xia Chen (Knoxville)

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be $p \ge 2$ independent identically distributed random walks started in the origin and taking values in \mathbb{Z}^d .

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The number of intersections of these walks can be measured in two natural ways: The intersection local time of the random walks,

$$I := \sum_{i_1=1}^{\infty} \cdots \sum_{i_p=1}^{\infty} \mathbf{1} \{ X^{(1)}(i_1) = \cdots = X^{(p)}(i_p) \},\$$

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counts the times when the paths intersect, whereas the intersection of the ranges

$$J:=\sum_{x\in\mathbb{Z}^d}\mathbf{1}\{X^{(1)}(i_1)=\cdots=X^{(p)}(i_p)=x ext{ for some } (i_1,\ldots,i_p)\}$$

counts the sites where the paths intersect.

Question: When are the random variables I and J finite?

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where $\mathbb{G}(x) \approx (|x|+1)^{2-d}$ is the Green's function.

Question: When are the random variables I and J finite?

Fact (Erdős and Taylor)

$$\mathbb{P}\{I < \infty\} = \mathbb{P}\{J < \infty\} = \begin{cases} 1 & \text{if } p(d-2) > d, \\ 0 & \text{otherwise.} \end{cases}$$

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From now on we assume that p(d-2) > d, i.e. we are in supercritical dimensions.

The tails of I and J

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The distributions of I and J are unknown. Khanin, Mazel, Shlosman, Sinai 1994 have studied the upper tail behaviour. They find that, for a sufficiently large,

$$\exp\big\{-c_1a^{\frac{1}{p}}\big\} \leq \mathbb{P}\big\{I > a\big\} \leq \exp\big\{-c_2a^{\frac{1}{p}}\big\}.$$

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Interestingly, the upper tails of J are substantially lighter. They show that, for all $\varepsilon > 0$ and all sufficiently large a,

$$\exp\left\{-a^{\frac{d-2}{d}+\varepsilon}\right\} \le \mathbb{P}\left\{J > a\right\} \le \exp\left\{-a^{\frac{d-2}{d}-\varepsilon}\right\}.$$

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The challenging question lies in understanding the difference of these behaviours, providing sharp estimates for the tails, and understanding the underlying 'optimal strategies' for the random walks.

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Let $W_2^{\varepsilon}(t)$ and $W_2^{\varepsilon}(t)$ be the ε -neighbourhoods of two independent Brownian paths starting at the origin and running for t time units.

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 $\lim_{t\uparrow\infty}\frac{1}{t^{(d-2)/d}}\log\mathbb{P}\big\{\big|W_1^\varepsilon(\theta t)\cap W_2^\varepsilon(\theta t)\big|\geq t\big\}=-I_d^\varepsilon(\theta),$

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and, if $d \ge 5$, there exists a critical θ^* such that

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This strongly suggests that, in the supercritical case $d \ge 5$,

$$\lim_{t\uparrow\infty}\frac{1}{t^{(d-2)/d}}\log\mathbb{P}\big\{\big|W_1^\varepsilon(\infty)\cap W_2^\varepsilon(\infty)\big|\geq t\big\}=-I_d^\varepsilon(\theta^*),$$

but their techniques do not allow the treatment of infinite times and this problem, like its discrete counterpart, remains open.

Let G be the Green's function of the random walk, defined by

$$\mathbb{G}(x) := \sum_{n=1}^{\infty} \mathbb{P}\{X(n) = x\}.$$

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Note that we are following the (slightly unusual) convention of not summing over the time n = 0, which has an influence on the value $\mathbb{G}(0)$.

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$$\mathfrak{A}_h \colon L^2(\mathbb{Z}^d) \to L^2(\mathbb{Z}^d)$$

is defined by

$$\mathfrak{A}_h g(x) = \sqrt{e^{h(x)} - 1} \sum_{y \in \mathbb{Z}^d} \mathbb{G}(x - y) g(y) \sqrt{e^{h(y)} - 1}.$$

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Our main result is formulated in terms of the spectral radius

$$\|\mathfrak{A}_{h}\|:=\sup\left\{\langle g,\mathfrak{A}_{h}g
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of the operator \mathfrak{A}_h .

Theorem 1 (Chen, M 2007)

The upper tail behaviour of the intersection local time I is given as

$$\lim_{a\uparrow\infty}\frac{1}{a^{1/p}}\log\mathbb{P}\big\{l>a\big\}=-p\,\inf\big\{\|h\|_q\,:\,h\ge 0\,\,\text{with}\,\,\|\mathfrak{A}_h\|\ge 1\big\}\,.$$

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Remark: The optimal strategy for the random walks is to each spend about $a^{1/p}$ time units in a bounded domain which does not grow with *a*. Then we get $I \approx a$ from intersections in this domain alone. This strategy makes *I* large without making *J* large, thus explaining the different tail behaviour.

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Selected ideas of the proof

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By a Tauberian theorem for any nonnegative X,

$$\lim_{k\uparrow\infty}\frac{1}{k}\log E\Big[\frac{X^k}{(k!)^p}\Big]=-\kappa\quad\Longleftrightarrow\quad\lim_{a\uparrow\infty}\frac{1}{a^{1/p}}\log P\{X>a\}=-pe^{\kappa/p}.$$

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$$I^{k} = \left[\sum_{x \in \mathbb{Z}^{d}} \prod_{j=1}^{p} \sum_{i=1}^{\infty} \mathbf{1}\{X^{(j)}(i) = x\}\right]^{k}$$

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$$= \sum_{x_{1},...,x_{k} \in \mathbb{Z}^{d}} \prod_{\ell=1}^{k} \prod_{j=1}^{p} \sum_{i=1}^{\infty} \mathbf{1}\{X^{(j)}(i) = x_{\ell}\}$$
$$= \sum_{x_{1},...,x_{k} \in \mathbb{Z}^{d}} \prod_{j=1}^{p} \sum_{i_{1},...,i_{k}=1}^{\infty} \prod_{\ell=1}^{k} \mathbf{1}\{X^{(j)}(i_{\ell}) = x_{\ell}\}$$

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$$\mathbb{E}I^k = \sum_{x_1,\ldots,x_k \in \mathbb{Z}^d} \prod_{j=1}^p \sum_{i_1,\ldots,i_k=1}^\infty \mathbb{E}\prod_{\ell=1}^k \mathbf{1}\{X^{(j)}(i_\ell) = x_\ell\}$$

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$$= \sum_{x_{1},\dots,x_{k}\in\mathbb{Z}^{d}}\left[\sum_{i_{1},\dots,i_{k}}\mathbb{E}\prod_{\ell=1}^{k}\mathbf{1}\{X(i_{\ell}) = x_{\ell}\}\right]^{p}$$

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$$\begin{split} \mathbb{E}I^{k} &= \sum_{x_{1},...,x_{k} \in \mathbb{Z}^{d}} \prod_{j=1}^{p} \sum_{i_{1},...,i_{k}=1}^{\infty} \mathbb{E} \prod_{\ell=1}^{k} \mathbf{1}\{X^{(j)}(i_{\ell}) = x_{\ell}\} \\ &= \sum_{x_{1},...,x_{k} \in \mathbb{Z}^{d}} \left[\sum_{i_{1},...,i_{k}} \mathbb{E} \prod_{\ell=1}^{k} \mathbf{1}\{X(i_{\ell}) = x_{\ell}\} \right]^{p} \\ &= \sum_{x_{1},...,x_{k} \in \mathbb{Z}^{d}} \left[\sum_{m=1}^{k} \sum_{\pi \in \mathcal{E}_{m}} \mathbf{1}\{(x_{1},...,x_{k}) \in \mathcal{A}(\pi)\} \sum_{\substack{j_{1},...,j_{m} \\ \text{distinct}}} \mathbb{E} \prod_{\ell=1}^{m} \mathbf{1}\{X(j_{\ell}) = x_{\pi_{\ell}}\} \right]^{p} \end{split}$$

where \mathcal{E}_m is the set of partitions $\{\pi_1, \ldots, \pi_m\}$ of $\{1, \ldots, k\}$ into *m* nonempty sets and $\mathcal{A}(\pi)$ is the set of tuples (x_1, \ldots, x_k) which are constant on the partitions.

$$\begin{split} \mathbb{E}I^{k} &= \sum_{x_{1},...,x_{k} \in \mathbb{Z}^{d}} \prod_{j=1}^{p} \sum_{i_{1},...,i_{k}=1}^{\infty} \mathbb{E}\prod_{\ell=1}^{k} \mathbf{1}\{X^{(j)}(i_{\ell}) = x_{\ell}\} \\ &= \sum_{x_{1},...,x_{k} \in \mathbb{Z}^{d}} \left[\sum_{i_{1},...,i_{k}} \mathbb{E}\prod_{\ell=1}^{k} \mathbf{1}\{X(i_{\ell}) = x_{\ell}\} \right]^{p} \\ &= \sum_{x_{1},...,x_{k} \in \mathbb{Z}^{d}} \left[\sum_{m=1}^{k} \sum_{\pi \in \mathcal{E}_{m}} \mathbf{1}\{(x_{1},...,x_{k}) \in \mathcal{A}(\pi)\} \sum_{\substack{j_{1},...,j_{m} \\ \text{distinct}}} \mathbb{E}\prod_{\ell=1}^{m} \mathbf{1}\{X(j_{\ell}) = x_{\pi_{\ell}}\} \right]^{p} \\ &= \sum_{x_{1},...,x_{k} \in \mathbb{Z}^{d}} \left[\sum_{m=1}^{k} \sum_{\pi \in \mathcal{E}_{m}} \mathbf{1}\{(x_{1},...,x_{k}) \in \mathcal{A}(\pi)\} \sum_{\sigma \in \mathfrak{S}_{m}} \prod_{\ell=1}^{m} \mathbb{G}\left(x_{\pi_{\sigma(\ell)}} - x_{\pi_{\sigma(\ell-1)}}\right) \right]^{p}, \end{split}$$

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Let $A \subset \mathbb{Z}^d$ be finite. Then we can analyse expressions of the form

$$\sum_{x_1,\ldots,x_k\in A} \left[\sum_{m=1}^k \sum_{\pi\in\mathcal{E}_m} \mathbf{1}\{(x_1,\ldots,x_k)\in\mathcal{A}(\pi)\}\sum_{\sigma\in\mathfrak{S}_m}\prod_{\ell=1}^m G(x_{\pi_{\sigma(\ell)}}-x_{\pi_{\sigma(\ell-1)}})\right]^p$$

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using

• large deviations for the empirical measure $L^{x} = \frac{1}{k} \sum_{j=1}^{k} \delta_{x_{j}}$,

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using

- large deviations for the empirical measure $L^{x} = \frac{1}{k} \sum_{j=1}^{k} \delta_{x_{j}}$,
- Hölder's inequality for the lower bound,
- the combinatorial fact that

$$\#\mathcal{E}_m = \frac{1}{m!} \sum_{\substack{j_1, \dots, j_m \ge 1\\ j_1 + \dots + j_m = k}} \frac{k!}{j_1! \cdots j_m!}$$

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$$\sum_{x_1,\ldots,x_k\in A} \left[\sum_{m=1}^k \sum_{\pi\in\mathcal{E}_m} \mathbf{1}\{(x_1,\ldots,x_k)\in\mathcal{A}(\pi)\}\sum_{\sigma\in\mathfrak{S}_m} \prod_{\ell=1}^m G(x_{\pi_{\sigma(\ell)}}-x_{\pi_{\sigma(\ell-1)}})\right]^p$$

using

- large deviations for the empirical measure $L^{x} = \frac{1}{k} \sum_{j=1}^{k} \delta_{x_{j}}$,
- Hölder's inequality for the lower bound,
- the combinatorial fact that

$$\#\mathcal{E}_{m} = \frac{1}{m!} \sum_{\substack{j_{1}, \dots, j_{m} \geq 1 \\ j_{1} + \dots + j_{m} = k}} \frac{k!}{j_{1}! \cdots j_{m}!}$$

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- an easy spectral theorem,
- symmetry (and nothing more!) of the function $G: \mathbb{Z}^d \to (0, \infty)$.

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We obtain, for finite $A \subset \mathbb{Z}^d$ that

$$\begin{split} \lim_{k \to \infty} \frac{1}{k} \log \frac{1}{k!} \sum_{x_1, \dots, x_k \in \mathcal{A}} \Big[\sum_{m=1}^k \sum_{\pi \in \mathcal{E}_m} \mathbf{1}_{\{(x_1, \dots, x_k) \in \mathcal{A}(\pi)\}} \sum_{\sigma \in \mathfrak{S}_m} \prod_{\ell=1}^m G(x_{\pi_{\sigma(\ell)}} - x_{\pi_{\sigma(\ell-1)}}) \Big]^p \\ &= -p \log \inf \Big\{ \|h\|_q \ : \ h \ge 0 \text{ with } \|\mathfrak{A}_h^A\| \ge 1 \Big\}, \end{split}$$

where the self-adjoint operator $\mathfrak{A}_{h}^{A} \colon L^{2}(A) \to L^{2}(A)$ is defined by

$$\mathfrak{A}_{h}^{A}g(x) = \sqrt{e^{h(x)} - 1} \sum_{y \in A} G(x - y) \sqrt{e^{h(y)} - 1} g(y).$$

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This suffices for the lower bound. The extension of the upper bound from finite sets A to the entire lattice is nontrivial, because the problem is not exponentially tight: Note that all shifts of A produce the same exponential decay of the upper tails of the intersection local times. To overcome this problem, we need to project the full problem onto a finite domain by wrapping it around a torus. The problem retains the given form, but with a different kernel G. We then let the period of the torus go to infinity.

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The upper tail behaviour of the intersection local time I is given as

$$\lim_{a\uparrow\infty}\frac{1}{a^{1/p}}\log\mathbb{P}\big\{l>a\big\}=-p\,\inf\big\{\|h\|_q\,:\,h\ge 0\,\,\text{with}\,\,\|\mathfrak{A}_h\|\ge 1\big\}\,.$$

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Remark: It is unsatisfactory that we cannot readily interpret the optimal h in the variational problem in a probabilistic manner. To some extent this is an artefact which is due to the discrete time structure of the random walk.

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For comparison we therefore now look at independent continuous time random walks

$$(X^{(1)}(t) \, : \, t \geq 0), \dots, (X^{(p)}(t) \, : \, t \geq 0)$$

and let A be their generator given by

$$Af(x) = \lim_{t\downarrow 0} \frac{\mathbb{E}_x f(X_t) - f(x)}{t}$$

A is a nonpositive definite, symmetric operator on $L^2(\mathbb{Z}^d)$.

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$$ilde{l} := \int_0^\infty dt_1 \cdots \int_0^\infty dt_p \ \mathbf{1} \{ X^{(1)}(t_1) = \cdots = X^{(p)}(t_p) \}.$$

Again we ask for the upper tail behaviour.

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Theorem 2 (Chen, M 2007)

The upper tail behaviour of the intersection local time \tilde{l} is given as

$$\lim_{a\uparrow\infty}\frac{1}{a^{1/\rho}}\log\mathbb{P}\big\{\tilde{I}>a\big\}=-\rho\,\inf\big\{\big\|\sqrt{-A}g\big\|_2^2\,\colon\,\|g\|_{2\rho}=1\big\}\,.$$

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Remark: The optimal strategy for the random walks is to have a local time field like

$$\ell^{(j)}(x) := \int_0^\infty \mathbf{1}\{X^{(j)}(t) = x\} \approx a^{1/p} g^2(x),$$

which implies

$$\tilde{l} = \sum_{x \in \mathbb{Z}^d} \prod_{j=1}^p \ell^{(j)}(x) \approx \sum_{x \in \mathbb{Z}^d} \prod_{j=1}^p a^{1/p} g^2(x) = a.$$

The probability of a random walk achieving such a local time is

$$pprox \exp\left\{-a^{1/p}\left\|\sqrt{-A}g\right\|_{2}^{2}
ight\},$$

which resembles the rate functions in Donsker-Varadhan theory.

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Our proof follows a similar strategy as in the discrete time case, but there is now an simpler formula for the kth moments

$$\mathbb{E}\tilde{I}^{k} = \sum_{x_{1},...,x_{k}\in\mathbb{Z}^{d}} \left[\sum_{\sigma\in\mathfrak{S}_{k}}\prod_{\ell=1}^{k}\mathbb{G}(x_{\sigma(\ell-1)}-x_{\sigma(\ell)})\right]^{p}.$$

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From this we obtain

$$\lim_{a\uparrow\infty}\frac{1}{a^{1/p}}\log\mathbb{P}\big\{\tilde{l}>a\big\}=-p\,\inf\big\{\|h\|_q\,:\,h\ge 0\,\,\text{with}\,\,\|\mathfrak{B}_h\|\ge 1\big\},$$

where the operator $\mathfrak{B}_h \colon L^2(\mathbb{Z}^d) \to L^2(\mathbb{Z}^d)$ is defined by

$$\mathfrak{B}_{h}g(x) = \sqrt{h(x)} \sum_{y \in \mathbb{Z}^{d}} \mathbb{G}(x-y)g(y) \sqrt{h(y)}$$

and the Green's function is

$$\mathbb{G}(x) = \int_0^\infty \mathbb{P}\{X(t) = x\} \, dt.$$

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where \mathfrak{G} is the Green's operator

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$$\mathfrak{G}f(x) := \sum_{y \in \mathbb{Z}^d} \mathbb{G}(x-y)f(y).$$

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The maximiser f exists and satisfies $\rho f = \mathfrak{G} f^{2p-1}$.

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The maximiser f exists and satisfies $\rho f = \mathfrak{G} f^{2p-1}$. We obtain the final form from $-A \circ \mathfrak{G} = id$ as

$$1/\rho = -\sup\left\{\langle f, Af \rangle \, : \, \|f\|_{2p} = 1\right\} = \inf\left\{\left\|\sqrt{-A}f\right\|_{2}^{2} \, : \, \|f\|_{2p} = 1\right\}.$$

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Concluding remarks

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We have obtained exact upper tail constants for the intersection local time of independent random walks in supercritical dimensions.

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- We have obtained exact upper tail constants for the intersection local time of independent random walks in supercritical dimensions.
- Our approach allows a direct treatment of the infinite time horizon avoiding the use of Donsker-Varadhan theory.
- We believe that this method has potential to solve some hard problems related to the intersection of the ranges as well. This work is in progress.

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