# Intersection of random walks in supercritical dimensions 

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joint work with Xia Chen (Knoxville)

## Framework

Let

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\left(X^{(1)}(n): n \in \mathbb{N}\right), \ldots,\left(X^{(p)}(n): n \in \mathbb{N}\right)
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The number of intersections of these walks can be measured in two natural ways: The intersection local time of the random walks,

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I:=\sum_{i_{1}=1}^{\infty} \cdots \sum_{i_{p}=1}^{\infty} 1\left\{X^{(1)}\left(i_{1}\right)=\cdots=X^{(p)}\left(i_{p}\right)\right\}
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counts the times when the paths intersect, whereas the intersection of the ranges

$$
J:=\sum_{x \in \mathbb{Z}^{d}} 1\left\{X^{(1)}\left(i_{1}\right)=\cdots=X^{(p)}\left(i_{p}\right)=x \text { for some }\left(i_{1}, \ldots, i_{p}\right)\right\}
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where $\mathbb{G}(x) \approx(|x|+1)^{2-d}$ is the Green's function.

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## Fact (Erdős and Taylor)

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\mathbb{P}\{I<\infty\}=\mathbb{P}\{J<\infty\}= \begin{cases}1 & \text { if } p(d-2)>d, \\ 0 & \text { otherwise. }\end{cases}
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From now on we assume that $p(d-2)>d$, i.e. we are in supercritical dimensions.

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Interestingly, the upper tails of $J$ are substantially lighter. They show that, for all $\varepsilon>0$ and all sufficiently large $a$,

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The challenging question lies in understanding the difference of these behaviours, providing sharp estimates for the tails, and understanding the underlying 'optimal strategies' for the random walks.

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and, if $d \geq 5$, there exists a critical $\theta^{*}$ such that

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This strongly suggests that, in the supercritical case $d \geq 5$,

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\lim _{t \uparrow \infty} \frac{1}{t^{(d-2) / d}} \log \mathbb{P}\left\{\left|W_{1}^{\varepsilon}(\infty) \cap W_{2}^{\varepsilon}(\infty)\right| \geq t\right\}=-I_{d}^{\varepsilon}\left(\theta^{*}\right)
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but their techniques do not allow the treatment of infinite times and this problem, like its discrete counterpart, remains open.

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Note that we are following the (slightly unusual) convention of not summing over the time $n=0$, which has an influence on the value $\mathbb{G}(0)$.

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Let $q>1$ be the conjugate of $p$, defined by $p^{-1}+q^{-1}=1$. For every nonnegative $h \in L^{q}\left(\mathbb{Z}^{d}\right)$ a bounded, symmetric, positive operator

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\mathfrak{A}_{h}: L^{2}\left(\mathbb{Z}^{d}\right) \rightarrow L^{2}\left(\mathbb{Z}^{d}\right)
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is defined by

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\mathfrak{A}_{h} g(x)=\sqrt{e^{h(x)}-1} \sum_{y \in \mathbb{Z}^{d}} \mathbb{G}(x-y) g(y) \sqrt{e^{h(y)}-1} .
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Our main result is formulated in terms of the spectral radius

$$
\left\|\mathfrak{A}_{n}\right\|:=\sup \left\{\left\langle g, \mathfrak{A}_{h} g\right\rangle:\|g\|_{2}=1\right\}
$$

of the operator $\mathfrak{A}_{h}$.

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Theorem 1 (Chen, M 2007)
The upper tail behaviour of the intersection local time $I$ is given as

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\lim _{a \nmid \infty \infty} \frac{1}{a^{1 / p}} \log \mathbb{P}\{I>a\}=-p \inf \left\{\|h\|_{G}: h \geq 0 \text { with }\left\|\mathscr{A}_{h}\right\| \geq 1\right\} .
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Remark: The optimal strategy for the random walks is to each spend about $a^{1 / p}$ time units in a bounded domain which does not grow with $a$. Then we get $I \approx a$ from intersections in this domain alone. This strategy makes I large without making $J$ large, thus explaining the different tail behaviour.

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By a Tauberian theorem for any nonnegative $X$,

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\lim _{k \uparrow \infty} \frac{1}{k} \log E\left[\frac{X^{k}}{(k!)^{p}}\right]=-\kappa \quad \Longleftrightarrow \quad \lim _{a \uparrow \infty} \frac{1}{a^{1 / p}} \log P\{X>a\}=-p e^{\kappa / p}
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& \mathbb{E} \ell^{k}=\sum_{X_{1}, \ldots, x \in \in Z^{j}} \prod_{j=1}^{p} \sum_{i, k, k=1}^{\infty} \mathbb{E} \prod_{\ell=1}^{k} 1\left\{X^{\omega}(i)=x_{k}\right\}
\end{aligned}
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where $\mathcal{E}_{m}$ is the set of partitions $\left\{\pi_{1}, \ldots, \pi_{m}\right\}$ of $\{1, \ldots, k\}$ into $m$ nonempty sets and $\mathcal{A}(\pi)$ is the set of tuples ( $x_{1}, \ldots, x_{k}$ ) which are constant on the partitions.

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& =\sum_{x_{1}, \ldots, x_{k} \in \mathbb{Z}^{d}}\left[\sum _ { m = 1 } ^ { k } \sum _ { \pi \in \mathcal { E } _ { m } } 1 \{ ( x _ { 1 } , \ldots , x _ { k } ) \in \mathcal { A } ( \pi ) \} \sum _ { \sigma \in \mathcal { S } _ { m } } \prod _ { \ell = 1 } ^ { m } \mathbb { T } _ { \ell = 1 } ^ { m } \left(x_{\pi_{\sigma(\ell)}}-x_{\pi},\right.\right.
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Let $A \subset \mathbb{Z}^{d}$ be finite. Then we can analyse expressions of the form

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\sum_{x_{1}, \ldots, x_{k} \in \mathcal{A}}\left[\sum_{m=1}^{k} \sum_{\pi \in \mathcal{E}_{m}} \mathbf{1}\left\{\left(x_{1}, \ldots, x_{k}\right) \in \mathcal{A}(\pi)\right\} \sum_{\sigma \in \mathcal{E}_{m}} \prod_{\ell=1}^{m} G\left(x_{\pi_{\sigma(\ell)}}-x_{\pi_{\sigma}(\ell-1)}\right)\right]^{p}
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- the combinatorial fact that

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\# \mathcal{E}_{m}=\frac{1}{m!} \sum_{\substack{j_{1}, \ldots, m_{m} \geq 1 \\ j_{1}+\ldots+t_{m}=k}} \frac{k!}{j_{1}!\cdots j_{m}!}
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\# \mathcal{E}_{m}=\frac{1}{m!} \sum_{\substack{j_{1}, \ldots, j_{m} \geq 1 \\ j_{1}+\cdots+t_{m}=k}} \frac{k!}{j_{1}!\cdots j_{m}!}
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- an easy spectral theorem,


## Selected ideas of the proof

Let $A \subset \mathbb{Z}^{d}$ be finite. Then we can analyse expressions of the form

$$
\sum_{x_{1}, \ldots, x_{k} \in A}\left[\sum_{m=1}^{k} \sum_{\pi \in \mathcal{E}_{m}} \mathbf{1}\left\{\left(x_{1}, \ldots, x_{k}\right) \in \mathcal{A}(\pi)\right\} \sum_{\sigma \in \mathfrak{S}_{m}} \prod_{\ell=1}^{m} G\left(x_{\pi_{\sigma(\ell)}}-x_{\pi_{\sigma(\ell-1)}}\right)\right]^{p}
$$

using

- large deviations for the empirical measure $L^{x}=\frac{1}{k} \sum_{j=1}^{k} \delta_{x_{j}}$,
- Hölder's inequality for the lower bound,
- the combinatorial fact that

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- an easy spectral theorem,
- symmetry (and nothing more!) of the function $G: \mathbb{Z}^{d} \rightarrow(0, \infty)$.


## Selected ideas of the proof

We obtain, for finite $A \subset \mathbb{Z}^{d}$ that

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\lim _{k \rightarrow \infty} \frac{1}{k} \log \frac{1}{k!} \sum_{x_{1}, \ldots, x_{k} \in A}\left[\sum_{m=1}^{k} \sum_{\pi \in \mathcal{E}_{m}} \mathbf{1}_{\left\{\left(x_{1}, \ldots, x_{k}\right) \in \mathcal{A}(\pi)\right\}} \sum_{\sigma \in \mathfrak{S}_{m}} \prod_{\ell=1}^{m} G\left(x_{\pi_{\sigma(\ell)}}-x_{\pi_{\sigma(\ell-1)}}\right)\right]^{p} \\
=-p \log \inf \left\{\|h\|_{q}: h \geq 0 \text { with }\left\|\mathfrak{A}_{h}^{A}\right\| \geq 1\right\}
\end{gathered}
$$

where the self-adjoint operator $\mathfrak{A}_{h}^{A}: L^{2}(A) \rightarrow L^{2}(A)$ is defined by

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\mathfrak{A}_{h}^{A} g(x)=\sqrt{e^{h(x)}-1} \sum_{y \in A} G(x-y) \sqrt{e^{h(y)}-1} g(y)
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This suffices for the lower bound. The extension of the upper bound from finite sets $A$ to the entire lattice is nontrivial, because the problem is not exponentially tight: Note that all shifts of $A$ produce the same exponential decay of the upper tails of the intersection local times. To overcome this problem, we need to project the full problem onto a finite domain by wrapping it around a torus. The problem retains the given form, but with a different kernel $G$. We then let the period of the torus go to infinity.

## Main result revisited

The upper tail behaviour of the intersection local time $I$ is given as

$$
\lim _{a \notinfty \infty} \frac{1}{a^{1 / p}} \log \mathbb{P}\{I>a\}=-p \inf \left\{\|h\|_{q}: h \geq 0 \text { with }\left\|\mathfrak{A}_{h}\right\| \geq 1\right\} .
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Remark: It is unsatisfactory that we cannot readily interpret the optimal $h$ in the variational problem in a probabilistic manner. To some extent this is an artefact which is due to the discrete time structure of the random walk.

A related problem

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For comparison we therefore now look at independent continuous time random walks

$$
\left(X^{(1)}(t): t \geq 0\right), \ldots,\left(X^{(p)}(t): t \geq 0\right)
$$

and let $A$ be their generator given by

$$
A f(x)=\lim _{t \downarrow 0} \frac{\mathbb{E}_{x} f\left(X_{t}\right)-f(x)}{t}
$$

$A$ is a nonpositive definite, symmetric operator on $L^{2}\left(\mathbb{Z}^{d}\right)$.

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Again we ask for the upper tail behaviour.

## A related problem

Theorem 2 (Chen, M 2007)
The upper tail behaviour of the intersection local time $\tilde{I}$ is given as

$$
\lim _{a \uparrow \infty} \frac{1}{a^{1 / p}} \log \mathbb{P}\{\tilde{l}>a\}=-p \inf \left\{\|\sqrt{-A} g\|_{2}^{2}:\|g\|_{2 p}=1\right\} .
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$$

Remark: The optimal strategy for the random walks is to have a local time field like

$$
\ell^{(j)}(x):=\int_{0}^{\infty} \mathbf{1}\left\{X^{(j)}(t)=x\right\} \approx a^{1 / p} g^{2}(x)
$$

which implies

$$
\tilde{I}=\sum_{x \in \mathbb{Z}^{d}} \prod_{j=1}^{p} \ell^{(j)}(x) \approx \sum_{x \in \mathbb{Z}^{d}} \prod_{j=1}^{p} a^{1 / p} g^{2}(x)=a .
$$

The probability of a random walk achieving such a local time is

$$
\approx \exp \left\{-a^{1 / p}\|\sqrt{-A} g\|_{2}^{2}\right\},
$$

which resembles the rate functions in Donsker-Varadhan theory.

## How do the limits compare?

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Our proof follows a similar strategy as in the discrete time case, but there is now an simpler formula for the $k$ th moments

$$
\mathbb{E} \tilde{I}^{k}=\sum_{x_{1}, \ldots, x_{k} \in \mathbb{Z}^{d}}\left[\sum_{\sigma \in \mathfrak{S}_{k}} \prod_{\ell=1}^{k} \mathbb{G}\left(x_{\sigma(\ell-1)}-x_{\sigma(\ell)}\right)\right]^{p}
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From this we obtain

$$
\lim _{a \uparrow \infty} \frac{1}{a^{1 / p}} \log \mathbb{P}\{\tilde{I}>a\}=-p \inf \left\{\|h\|_{q}: h \geq 0 \text { with }\left\|\mathfrak{B}_{h}\right\| \geq 1\right\}
$$

where the operator $\mathfrak{B}_{h}: L^{2}\left(\mathbb{Z}^{d}\right) \rightarrow L^{2}\left(\mathbb{Z}^{d}\right)$ is defined by

$$
\mathfrak{B}_{h} g(x)=\sqrt{h(x)} \sum_{y \in \mathbb{Z}^{d}} \mathbb{G}(x-y) g(y) \sqrt{h(y)}
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and the Green's function is

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\mathbb{G}(x)=\int_{0}^{\infty} \mathbb{P}\{X(t)=x\} d t
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& \quad=\inf \left\{b: \sup _{\substack{\|g\|_{2}=1 \\
\|h\|_{q}=1}}\langle\sqrt{h} g, \mathfrak{G} \sqrt{h} g\rangle \geq 1 / b\right\}
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\mathfrak{G} f(x):=\sum_{y \in \mathbb{Z}^{d}} \mathbb{G}(x-y) f(y)
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The maximiser $f$ exists and satisfies $\rho f=\mathfrak{G} f^{2 p-1}$. We obtain the final form from $-A \circ \mathfrak{G}=i d$ as

$$
1 / \rho=-\sup \left\{\langle f, A f\rangle:\|f\|_{2 p}=1\right\}=\inf \left\{\|\sqrt{-A} f\|_{2}^{2}:\|f\|_{2 p}=1\right\}
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- We have obtained exact upper tail constants for the intersection local time of independent random walks in supercritical dimensions.
- Our approach allows a direct treatment of the infinite time horizon avoiding the use of Donsker-Varadhan theory.
- We believe that this method has potential to solve some hard problems related to the intersection of the ranges as well. This work is in progress.

