

Coercivity of high-frequency scattering problems

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Joint work with: Euan Spence (Bath),
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Spence, Chandler-Wilde, Graham, S, Comm Pure Appl Math 2011.

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Summary

- ▶ ~ 1960's - 1980's: tremendous interest in rigorous aspects of scattering/ diffraction:

- ▶ decay at $t \rightarrow \infty$ of wave equation $\frac{\partial^2 w}{\partial t^2} - c^2 \Delta w = 0$

- ▶ asymptotics as $k \rightarrow \infty$ of Helmholtz equation $\Delta u + k^2 u = 0$

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 - ▶ asymptotics as $k \rightarrow \infty$ of Helmholtz equation $\Delta u + k^2 u = 0$
 - ▶ (related in subtle way)
- ▶ ~ 2000's - present: interest in Numerical Analysis of Helmholtz for $k \gg 1$
 - ▶ e.g. "hybrid asymptotic-numerical methods"
 - ▶ This talk: boundary integral equations
 - ▶ One analysis question: prove relevant operator is **coercive**
 - ▶ Surprise (?) – appears that for coercivity need **stronger** results than those obtained classically (at least in the context of one classic tool – **Morawetz multipliers**)

Obstacle scattering problem: acoustically soft/ TE Maxwell (2d) perfectly conducting boundary

$$u^i = e^{ikx \cdot \hat{a}}$$

$$u^s$$

Ω_{int}

Γ

$$\Delta u^s + k^2 u^s = 0$$

$$\text{in } \Omega_{\text{ext}} := \mathbb{R}^d \setminus \Omega_{\text{int}}$$

$$u^s + u^i = 0 \text{ on } \Gamma$$

$$k > 0$$

Radiation conditions:

$$\frac{\partial u^s}{\partial r} - iku^s = o\left(r^{-\frac{d-1}{2}}\right) \text{ as } r \rightarrow \infty.$$

\implies Uniqueness and existence.

Boundary integral equations

- ▶ Green's Integral Representation:

$$u^s(x) = \int_{\Gamma} \left(\frac{\partial \Phi_k}{\partial n(y)}(x, y) u^s(y) - \Phi_k(x, y) \frac{\partial u^s}{\partial n}(y) \right) ds(y), \quad x \in \Omega_{\text{ext}}$$

where

$$\Phi_k(x, y) := \begin{cases} \frac{i}{4} H_0^{(1)}(k|x-y|) & (d=2) \\ \frac{e^{ik|x-y|}}{4\pi|x-y|} & (d=3) \end{cases}$$

Hence boundary integral equations for $v := \frac{\partial u}{\partial n}$:

- ▶ Single layer:

$$S_k v(x) := \int_{\Gamma} \Phi_k(x, y) v(y) ds(y) = u^i(x)$$

(uniqueness fails for $k^2 =$ interior Dirichlet eigenvalues)

- ▶ (Adjoint) double layer:

$$\left(\frac{1}{2} I + D'_k \right) v(x) := \frac{1}{2} v(x) + \int_{\Gamma} \frac{\partial \Phi_k}{\partial n(x)}(x, y) v(y) ds(y) = \frac{\partial u^i}{\partial n}(x)$$

(uniqueness fails for $k^2 =$ interior Neumann eigenvalues)

Combined boundary integral equations

Try a *combination* of a double layer and of a single layer:

$$(\text{Double Layer}) - i\eta \times (\text{Single Layer})$$

with a 'coupling constant' $\eta \sim k$ ($k \gg 1$).

i.e. let

$$A_k = \frac{1}{2}I + D'_k - i\eta S_k.$$

- ▶ \rightsquigarrow 'Combined' boundary integral equation:

$$A_k \left(\frac{\partial u}{\partial n} \right) = f \quad \left(f = \frac{\partial u^i}{\partial n}(x) - i\eta u^i(x) \right)$$

- ▶ At *high frequencies* ($k \gg 1$) kernel of A_k highly oscillatory (and non-linear) in k .

The operator A_k

$$A_k \left(\frac{\partial u}{\partial n} \right) = f$$

- ▶ For a fixed k , $\eta > 0$: $A_k : L^2(\Gamma) \rightarrow L^2(\Gamma)$ bounded and **invertible** and

$$\partial u / \partial n \in L^2(\Gamma) \text{ if } \Gamma \text{ is Lipschitz (Nečas)}$$

- ▶ Q. What do we want to know about A_k ?

1. bound on $\|A_k\|$ (**explicit in k**)
2. bound on $\|A_k^{-1}\|$ (**explicit in k**)
3. **coercivity**: $\exists \gamma > 0$ such that

$$|(A_k \phi, \phi)_{L^2(\Gamma)}| \geq \gamma \|\phi\|_{L^2(\Gamma)}^2, \quad \forall \phi \in L^2(\Gamma)$$

(**γ explicit in k**)

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- ▶ Q. What do we want to know about A_k ?
 1. bound on $\|A_k\|$ (**explicit in k**) ← **relatively easy**
 2. bound on $\|A_k^{-1}\|$ (**explicit in k**)
 3. **coercivity**: $\exists \gamma > 0$ such that

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 1. bound on $\|A_k\|$ (**explicit in k**) ← relatively easy
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(**γ explicit in k**) ← even harder!

Why is bounding $\|A_k^{-1}\|$ not enough?

$$A_k v = f \quad \left(v = \frac{\partial u}{\partial n} \right)$$

- ▶ **Solving numerically using Galerkin method:** choose $\mathcal{S}_N \subset L^2(\Gamma)$ (N -dimensional subspace), find $v_N \in \mathcal{S}_N$ such that

$$(A_k v_N, \phi_N)_{L^2(\Gamma)} = (f, \phi_N)_{L^2(\Gamma)}, \quad \forall \phi_N \in \mathcal{S}_N$$

- ▶ Want “quasi-optimality”: (Lax-Milgram + Cea’s Lemma) :

$$\|v - v_N\|_{L^2(\Gamma)} \leq C(k) \inf_{\phi_N \in \mathcal{S}_N} \|v - \phi_N\|_{L^2(\Gamma)} \quad (*)$$

–in some sense “numerical well-posedness”

- ▶ $C(k) = \|A_k\|/\gamma_k \therefore$ **Bound on $\|A_k^{-1}\|$ can’t give (*) for important \mathcal{S}_N**

Plan

Multiplier Methods

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Helmholtz equation

$$\Delta u + k^2 u = 0,$$

where $u(x)$, $x \in D \subset \mathbb{R}^3$, $k > 0$

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Helmholtz equation

$$\int_D \bar{M} \left(\Delta u + k^2 u \right) = 0,$$

where $u(x)$, $x \in D \subset \mathbb{R}^3$, $k > 0$

integrate by parts

$$\bar{M} \Delta u = \nabla \cdot (\bar{M} \nabla u) - \nabla \bar{M} \cdot \nabla u$$

get

$$\int_{\partial D} \bar{M} \frac{\partial u}{\partial n} - \int_D \nabla \bar{M} \cdot \nabla u + k^2 \int_D \bar{M} u = 0$$

Some famous (and not so famous) multipliers

$$\int_D \overline{M} \left(\Delta u^s + k^2 u^s \right) = 0,$$

radiation condition for u^s :

$$\therefore \frac{\partial u^s}{\partial r} - iku^s = o\left(r^{-\frac{d-1}{2}}\right), \quad d = 2, 3$$

$$u^s(x) \sim \frac{e^{ikr}}{r^{\frac{d-1}{2}}} f(\hat{x}) \quad \text{as } r = |x| \rightarrow \infty$$

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- ▶ Green (1828), $M = u^s$
- ▶ Rellich (1940), e.g. $M = r \frac{\partial u^s}{\partial r} = x \cdot \nabla u^s$
- ▶ Morawetz (1968), e.g. $M = r \frac{\partial u^s}{\partial r} - ikr u^s + \frac{d-1}{2} u^s$

Classic (high frequency) scattering/ diffraction theory

- ▶ Enormous interest from 1960's onwards, e.g.,
 - ▶ USA – Keller, Lax, Philips, Morawetz (@ Courant), Melrose...
 - ▶ USSR/ Russia - Fock, Buslaev, Babich, Vainberg...
- ▶ 3 main problems
 1. Wave equation: behaviour as $t \rightarrow \infty$
 2. Wave equation: propagation of singularities
 3. Helmholtz: behaviour as $k \rightarrow \infty$
- ▶ related in subtle way: “1+2=3” [Vainberg, 1975]

Key concept: (non-)trapping

- ▶ as $k \rightarrow \infty$ Helmholtz in trapping domains has “almost eigenvalues/eigenfunctions” (resonances)

Key concept: (non-)trapping

- ▶ as $k \rightarrow \infty$ Helmholtz in trapping domains has “almost eigenvalues/eigenfunctions” (resonances)
- ▶ Classic theory can be translated into results about A_k^{-1} .
- ▶ Expect that

1. For Ω_{ext} certain *trapping* domains

$$\|A_{k_n}^{-1}\| \gtrsim e^{\alpha k_n}, \quad 0 < k_1 < k_2 < \dots \text{ some } \alpha > 0$$

2. If Ω_{ext} is *non-trapping* then

$$\|A_k^{-1}\| \lesssim 1, \quad \forall k \geq k_0$$

▶ Find

1. Proved
2. Proved for star-shaped domains [Chandler-Wilde, Monk 2008] using Rellich $\left(M = \frac{\partial u}{\partial r}\right)$ (N.B. needed extra work to deal with ∞)

The operator A_k

$$A_k \left(\frac{\partial u}{\partial n} \right) = f$$

- ▶ Spaces: $A_k : L^2(\Gamma) \rightarrow L^2(\Gamma)$ bounded and invertible and $\partial u / \partial n \in L^2(\Gamma)$ if Γ is Lipschitz
- ▶ Q. What do we want to know about A_k ?
 1. bound on $\|A_k\|$ (explicit in k) ← relatively easy
 2. bound on $\|A_k^{-1}\|$ (explicit in k) ← use classic high-frequency scattering theory
 3. coercivity: $\exists \gamma > 0$ such that

$$|(A_k \phi, \phi)_{L^2(\Gamma)}| \geq \gamma \|\phi\|_{L^2(\Gamma)}^2, \quad \forall \phi \in L^2(\Gamma)$$

(γ explicit in k) ← why?

Quasi-optimality

$$\|v - v_N\|_{L^2(\Gamma)} \leq C(k) \inf_{\phi_N \in \mathcal{S}_N} \|v - \phi_N\|_{L^2(\Gamma)} \quad (\star)$$

- ▶ Want to establish (with explicit k dependence of C) for
 1. \mathcal{S}_N piecewise polynomials
 2. $\mathcal{S}_{N,k}$ “hybrid” subspace incorporating asymptotics of $v = \frac{\partial u^s}{\partial n}$
- ▶ For 1. need $N = \mathcal{O}(k^{d-1})$ as $k \rightarrow \infty$, possibility of 2. giving $N = \mathcal{O}(1)$.
- ▶ k -explicit (\star) for 1. – classic problem, solved by [Melenk, 2011] (needs bound on $\|A_k^{-1}\|$)
- ▶ Coercivity (+bound on $\|A_k\|$ – easy) gives (\star) for 1. and 2. k -explicit.

Coercivity

$\exists \gamma > 0$ such that

$$|(A_k \phi, \phi)_{L^2(\Gamma)}| \geq \gamma \|\phi\|_{L^2(\Gamma)}^2, \quad \forall \phi \in L^2(\Gamma)$$

- ▶ Not obvious will hold – standard approach to formulations of Helmholtz: prove

$$\text{operator}_k = \text{coercive} + \text{compact}_k$$

- ▶ Coercivity for circle (2d) and sphere (3d) $\forall k \geq k_0$, $\gamma = 1$
[Domínguez, Graham, Smyshlyaev, 2007] (Fourier analysis)

Two Coercivity Results using Morawetz Multipliers

Result 1. Ω_{int} Lipschitz star-shaped, a specially constructed “star-combined” \mathcal{A}_k is coercive $\forall k, \gamma = \mathcal{O}(1)$. [Spence, Chandler-Wilde, Graham, S., *Comm Pure Appl Math* 2011]

Result 2. Ω_{int} smooth convex, the classical combined A_k is coercive $\forall k \geq k_0, \eta > \eta_0 k, \gamma = \frac{1}{2} - \varepsilon$ [Spence, I.Kamotski, S. *Comm Pure Appl Math* 2015]

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How?

- ▶ A_k arose from $\frac{\partial u^s}{\partial n} - i\eta u^s$ on Γ
- ▶ $M = r \frac{\partial u^s}{\partial r} - ikr u^s + \frac{d-1}{2} u^s \rightsquigarrow$ star-shaped coercivity 1.
- ▶ $M = Z(x) \cdot \nabla u^s - i\eta(x) u^s + \alpha(x) u^s \rightsquigarrow$ smooth convex coercivity 2.

Morawetz - 1

Morawetz & Ludwig (1968): take as a multiplier

$$Mu := \mathbf{x} \cdot \nabla u - ikru + \frac{d-1}{2}u,$$

Then (the Morawetz-Ludwig identity):

$$2\operatorname{Re}(\overline{Mu}(\Delta u + k^2u)) =$$

$\nabla \cdot \left[2\operatorname{Re}(\overline{Mu}\nabla u) + (k^2|u|^2 - |\nabla u|^2)\mathbf{x} \right] - \left(|\nabla u|^2 - |u_r|^2 \right) - |u_r - iku|^2$. A corollary:

$$\int_{\partial D} \left[2\operatorname{Re} \left(\overline{Mu} \frac{\partial u}{\partial \nu} \right) + (k^2|u|^2 - |\nabla u|^2)\mathbf{x} \cdot \nu \right] ds =$$

$$\int_D \left[2\operatorname{Re}(\overline{Mu}(\Delta u + k^2u)) + \left(|\nabla u|^2 - |u_r|^2 \right) + |u_r - iku|^2 \right] dx$$

Let Ω_{in} star-shaped $\Leftrightarrow \mathbf{x} \cdot \mathbf{n} \geq \beta > 0$ ($\mathbf{n} = -\nu$). Take

$D = \Omega_{ext} \cap B(0, R)$, let $R \rightarrow \infty$, use radiation conditions. \implies k -uniform bounds for D-t-N; error bounds for GO/ GTD, etc.

Morawetz - 1

“Star-combined” BIE (Spence, Chandler-Wilde, Graham, S., *Comm Pure Appl Math* 2011):

In the Morawetz identity, choose $u = S_k \phi$. Take $D = (\Omega_{\text{ext}} \cap B(0, R)) \cup \Omega_{\text{int}}$, let $R \rightarrow \infty$. Then, for

$$\mathcal{A}_k := (\mathbf{x} \cdot \mathbf{n}) \left(\frac{1}{2} I + D'_k \right) + \mathbf{x} \cdot \nabla_{\Gamma} S_k - i \eta S_k, \quad \eta := kr + i \frac{d-1}{2}.$$

$$\mathcal{A}_k \left(\frac{\partial u}{\partial n} \right) = f \quad (f = \mathbf{x} \cdot \nabla u^i(x) - i \eta(x) u^i(x))$$

Coercivity: $\forall k \geq 0$,

$$\operatorname{Re}(\mathcal{A}_k \phi, \phi) \geq \gamma \|\phi\|_{L^2}^2, \quad \gamma = \frac{1}{2} \operatorname{ess\,inf}_{x \in \Gamma} (\mathbf{x} \cdot \mathbf{n}(x)) > 0.$$

Classical combined (Spence, I. Kamotski, S., *Comm Pure Appl Math* 2015)

Try as a multiplier, with appropriate vector field $\mathbf{Z}(\mathbf{x})$, and scalar functions $\alpha(\mathbf{x})$ and $\beta(\mathbf{x})$:

$$Mu := \mathbf{Z}(\mathbf{x}) \cdot \nabla u - ik\beta(\mathbf{x})u + \alpha(\mathbf{x})u.$$

Then the following Morawetz-type identity holds:

$$\begin{aligned} 2\operatorname{Re}(\overline{Mu}(\Delta + k^2)u) = \\ \nabla \cdot \left[2\operatorname{Re}(\overline{Mu} \nabla u) + (k^2|u|^2 - |\nabla u|^2)\mathbf{Z} \right] + \\ (2\alpha - \nabla \cdot \mathbf{Z})(k^2|u|^2 - |\nabla u|^2) \\ - 2\operatorname{Re}(\partial_i Z_j \partial_i u \overline{\partial_j u}) - 2\operatorname{Re}(\overline{u}(ik\nabla\beta + \nabla\alpha) \cdot \nabla u) \end{aligned}$$

For wave equation Morawetz needed:

- ▶ $Z(x), \quad x \in \Omega_{\text{ext}}$
- ▶ $\Re(\partial_j Z_i \xi_i \bar{\xi}_j) \geq 0, \quad \xi \in \mathbb{C}^d,$
- ▶ $Z.n > 0$ on $\Gamma,$
- ▶ $Z(x) \rightarrow c x$ as $|x| \rightarrow \infty$

(almost enough for $\|A_k^{-1}\|$ bound)

have for Ω_{ext} non-trapping
in 2-d

For coercivity of A_k we need:

- ▶ $Z(x), \quad x \in \Omega_{\text{ext}} \cup \Omega_{\text{int}}$
- ▶ $\Re(\partial_j Z_i \xi_i \bar{\xi}_j) \geq \theta |\xi|^2, \quad \xi \in \mathbb{C}^d,$
- ▶ $Z = n$ on $\Gamma,$
- ▶ $Z(x) \rightarrow c x$ as $|x| \rightarrow \infty$

have for Ω_{int} smooth convex
in 2 & 3-d

...non-trapping?

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