

# Introduction to the gradient flow in QCD

Susanne Ehret

Higgs Centre for Theoretical Physics, University of Edinburgh



YTF - December 2014 - Durham

# Content

- Motivation
- Gradient flow in QCD
- $D+1$  dimensional theory and renormalisation
- Energy momentum tensor
- Outlook

# Motivation

- Bare, local, composite operator

$$O_R = \sum_P Z_{OP} (P_0 - \langle P_0 \rangle)$$

- Leads to problems with cancellation of 'infinities' in lattice calculations

# Motivation

- Bare, local, composite operator

$$O_R = \sum_P Z_{OP} (P_0 - \langle P_0 \rangle)$$

- Leads to problems with cancellation of 'infinities' in lattice calculations
- New method: gradient flow

# Motivation

- Bare, local, composite operator

$$O_R = \sum_P Z_{OP} (P_0 - \langle P_0 \rangle)$$

- Leads to problems with cancellation of 'infinities' in lattice calculations
- New method: gradient flow
- Define energy momentum tensor in a meaningful way on the lattice to study e.g. scaling behaviour of QFTs

# Motivation

- Bare, local, composite operator

$$O_R = \sum_P Z_{OP} (P_0 - \langle P_0 \rangle)$$

- Leads to problems with cancellation of 'infinities' in lattice calculations
- New method: gradient flow
- Define energy momentum tensor in a meaningful way on the lattice to study e.g. scaling behaviour of QFTs
- Dunne, Resurgence: Gradient flow for sign problem, steepest decent

# The gradient flow in pure gauge theory

Euclidean action

$$S = -\frac{1}{2} \int d^D x \operatorname{tr} \{ F_{\mu\nu}^2 \}$$

# The gradient flow in pure gauge theory

Euclidean action

$$S = -\frac{1}{2} \int d^D x \operatorname{tr} \{ F_{\mu\nu}^2 \}$$

Gradient flow

$$\partial_t B_\mu(t, x) = -\frac{\delta S}{\delta B_\mu}, \quad B_\mu(t, x)|_{t=0} = A_\mu(x)$$



# The gradient flow in pure gauge theory

Euclidean action

$$S = -\frac{1}{2} \int d^D x \operatorname{tr} \{F_{\mu\nu}^2\}$$

Gradient flow

$$\partial_t B_\mu(t, x) = -\frac{\delta S}{\delta B_\mu}, \quad B_\mu(t, x)|_{t=0} = A_\mu(x)$$

Flow equation

$$\partial_t B_\mu(t, x) = D_\nu G_{\nu\mu}$$

$$D_\nu = \partial_\nu + g_0 B_\nu, \quad G_{\nu\mu} = \frac{1}{g_0} [D_\nu, D_\mu]$$

# The gradient flow in pure gauge theory - smearing

Heat equation

$$\partial_t B_\mu = \partial^2 B_\mu + O(g_0)$$

$$B_\mu|_{t=0} = A_\mu$$

# The gradient flow in pure gauge theory - smearing

Heat equation

$$\partial_t B_\mu = \partial^2 B_\mu + O(g_0)$$

$$B_\mu|_{t=0} = A_\mu$$

Solution: heat kernel

$$B_\mu(t, x) = \int d^D y K_t(x - y)_{\mu\nu} A_\nu(y) + O(g_0)$$

$$K_t(x) = \frac{e^{-\frac{x^2}{4t}}}{(4\pi t)^{D/2}}$$

# The gradient flow in pure gauge theory - smearing

Heat equation

$$\partial_t B_\mu = \partial^2 B_\mu + O(g_0)$$

$$B_\mu|_{t=0} = A_\mu$$

Solution: heat kernel

$$B_\mu(t, x) = \int d^D y K_t(x - y)_{\mu\nu} A_\nu(y) + O(g_0)$$

$$K_t(x) = \frac{e^{-\frac{x^2}{4t}}}{(4\pi t)^{D/2}}$$

Exponential damping, smearing with radius  $\sqrt{8t}$

# Field Theory in $D+1$ dimensions and additional Feynman rules

$$S_{tot} = S + S_{fl}$$

Feynman rules

# Field Theory in $D+1$ dimensions and additional Feynman rules

$$S_{tot} = S + S_{fl} + S_{gf} + S_{c\bar{c}}$$

Feynman rules

# Field Theory in $D+1$ dimensions and additional Feynman rules

$$S_{tot} = S + S_{fl} + S_{gf} + S_{c\bar{c}}$$

$$S_{fl} = \int_0^\infty dt \int d^D x L_\mu(t, x) (\partial_t B_\mu - D_\nu G_{\nu\mu})(t, x)$$

Feynman rules

# Field Theory in $D+1$ dimensions and additional Feynman rules

$$S_{tot} = S + S_{fl} + S_{gf} + S_{c\bar{c}}$$

$$S_{fl} = \int_0^\infty dt \int d^D x L_\mu(t, x) (\partial_t B_\mu - D_\nu G_{\nu\mu})(t, x)$$

Feynman rules



# Field Theory in $D+1$ dimensions and additional Feynman rules

$$S_{tot} = S + S_{fl} + S_{gf} + S_{c\bar{c}}$$

$$S_{fl} = \int_0^\infty dt \int d^D x L_\mu(t, x) (\partial_t B_\mu - D_\nu G_{\nu\mu})(t, x)$$

Feynman rules

$$BB \propto \frac{1}{p^2} e^{-tp^2}$$

$$X^{(2/3)} \propto V^{(3/4)}$$

# Field Theory in $D+1$ dimensions and additional Feynman rules

$$S_{tot} = S + S_{fl} + S_{gf} + S_{c\bar{c}}$$

$$S_{fl} = \int_0^\infty dt \int d^D x L_\mu(t, x) (\partial_t B_\mu - D_\nu G_{\nu\mu})(t, x)$$

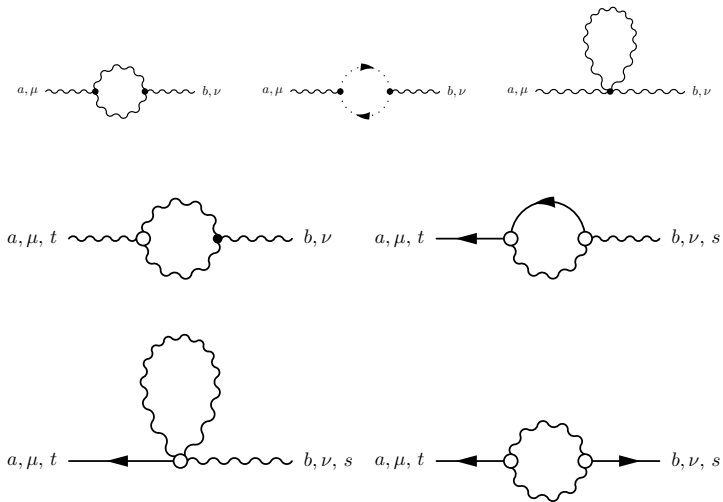
Feynman rules

$$BB \propto \frac{1}{p^2} e^{-tp^2}$$

$$X^{(2/3)} \propto V^{(3/4)}$$

$$BL \propto \theta(t-s) \tilde{K}_t(p) \propto \theta(t-s) e^{-tp^2}$$

# Pure gauge



# Renormalisation of the bulk field

- Pure gauge: no additional renormalisation for bulk field required
- Prove to all orders with BRST (Lüscher)
- Special for gauge case; not trivial for fermions
- In general: divergencies  $\propto \delta(t) \Rightarrow$  only at the boundary
- Seen from  $e^{-tp^2}$ : large momenta suppressed
- Fermions: similar treatment, multiplicative renormalisation

$$\chi = Z_\chi^{-1/2} \chi_R, \quad L = Z_\chi^{1/2} L_R$$

# Energy momentum tensor

- Interesting for many reasons:
  - Equation of state
  - Quark gluon plasma
  - BSM: Scale invariance, conformal window

# Energy momentum tensor

- Interesting for many reasons:
  - Equation of state
  - Quark gluon plasma
  - BSM: Scale invariance, conformal window
- Non-perturbative: lattice
- But: no translational invariance: need definition that recovers scale invariance in continuum limit
- EMT not finite anymore (Ward identity)
- Problem with subtraction of "infinities"
- Gradient flow might solve this last problem

## Summary and Outlook

- Gauge fields at non-zero flow time need no additional renormalisation
- Fermions receive additional bulk field renormalisation

## Summary and Outlook

- Gauge fields at non-zero flow time need no additional renormalisation
- Fermions receive additional bulk field renormalisation
- Define renormalisation independent observables and finite composite operators
- General aspects of non-Abelian field theories, e.g. scale invariance due to a well-defined EMT on lattice
- High precision measurements



## Summary and Outlook

- Gauge fields at non-zero flow time need no additional renormalisation
- Fermions receive additional bulk field renormalisation
- Define renormalisation independent observables and finite composite operators
- General aspects of non-Abelian field theories, e.g. scale invariance due to a well-defined EMT on lattice
- High precision measurements

Thank you!