

Constraining simplified dark matter models with the LHC

Karl Nordström¹ and Thomas Jacques²

¹University of Glasgow

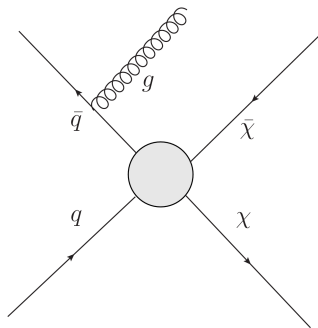
²Université de Genève

December 18th, 2014

Some brief background:

- ▶ The LHC can investigate dark matter (DM) models where there is some way for the dark sector to talk to light quarks
- ▶ Can roughly separate searches into two types:
 1. Model-dependent (SUSY searches, ...)
 2. Model-agnostic (typically mono-X)

I will discuss the models used for setting limits on 2. and present some constraints using monojet (a single jet + E_T^{miss}) limits in particular



We will assume:

- ▶ Dirac fermion dark matter χ with mass m_{DM}
- ▶ A vector mediator Z' with mass M and pure axial-vector¹ couplings g_q, g_{DM}

$$\Rightarrow \mathcal{L}_{\text{MSDM}} \supset - \sum_q g_q Z'_\mu \bar{q} \gamma^\mu \gamma_5 q - g_{\text{DM}} Z'_\mu \bar{\chi} \gamma^\mu \gamma_5 \chi \quad (1)$$

This is the interaction term of our minimal simplified dark matter model (MSDM). There are four free parameters ($M, m_{\text{DM}}, g_q, g_{\text{DM}}$).

¹The LHC has little sensitivity to vector couplings compared to direct detection (spin-independent vs spin-dependent).

Can expand the mediator propagator:

$$\frac{g_q g_{\text{DM}}}{Q^2 - M^2} \sim -\frac{g_q g_{\text{DM}}}{M^2} \left(1 + \frac{Q^2}{M^2} + \mathcal{O}\left(\frac{Q^4}{M^4}\right) \right) \quad (2)$$

Let $\Lambda = M/\sqrt{g_q g_{\text{DM}}}$, then:

$$\mathcal{L}_{\text{EFT}} \supset \frac{1}{\Lambda^2} \chi \gamma^\mu \gamma_5 \bar{\chi} q \gamma_\mu \gamma_5 \bar{q} \quad (3)$$

where we have integrated out the mediator using (2). This is an effective interaction term which is valid when $Q \ll M$. We've reduced the number of free parameters to two (Λ , m_{DM}).

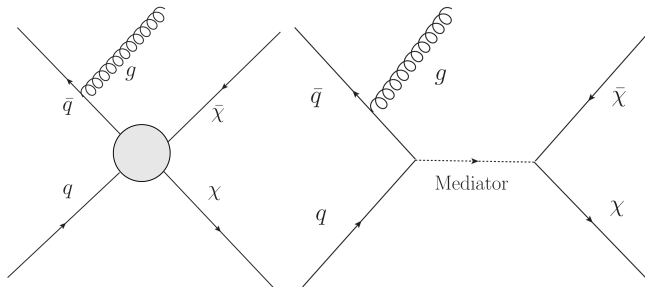


Figure : Example of what we are working with.

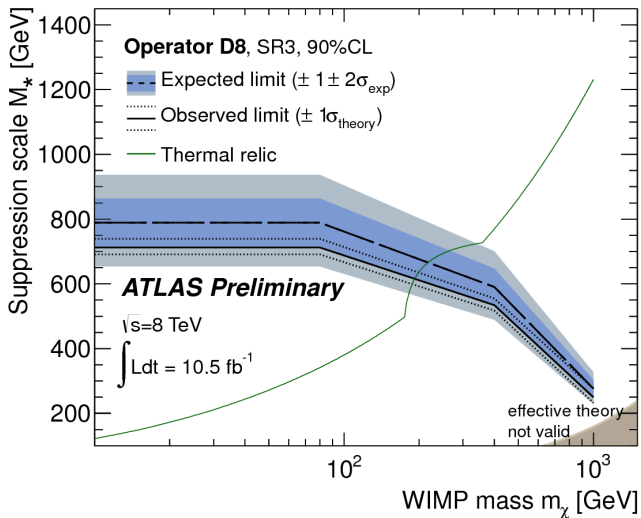


Figure : Example of ATLAS limits on our EFT operator.

To facilitate the comparison to direct detection constraints² LHC experiments have generally interpreted model-agnostic searches using EFTs.

But:

- ▶ Direct detection: $Q \sim \mathcal{O}(10 \text{ keV}) \Rightarrow$ EFT valid \sim always
- ▶ LHC: $Q \sim \mathcal{O}(1 \text{ TeV}) \Rightarrow$ EFT valid \sim ?

Answer (1307.2253, 1308.6799 + others): EFT is only valid for $M \gtrsim 2.5 \text{ TeV}$ at $\sqrt{s} = 8 \text{ TeV}$.

²And also just because there are fewer parameters \Rightarrow cheaper, easier.

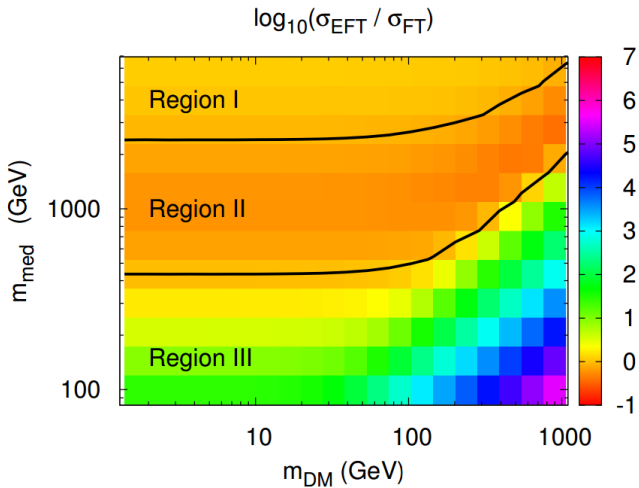


Figure : Ratio of simplified model to EFT cross-section for $g_q, g_{\text{DM}} = 1$ (from 1308.6799).

- ▶ Some recent studies using the CMS limits and NLOPS predictions for $\chi\bar{\chi} + 1$ jet (1407.8257, 1411.0535)
- ▶ Generally present constrains in e.g. the $M - m_{\text{DM}}$ plane as an exclusion contour for a particular choice of g_q, g_{DM}
- ▶ Note that the minimum width Γ_{min} can be calculated from the input parameters and needs to be taken into account!

We have studied constraints using ATLAS limits and LOPS predictions scanning g_q/g_{DM} and $g_q \cdot g_{\text{DM}}$ assuming $\Gamma_M = \Gamma_{\text{min}}$.

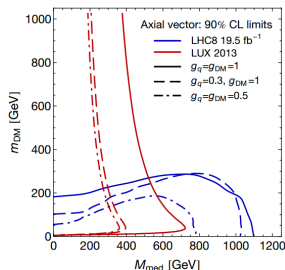
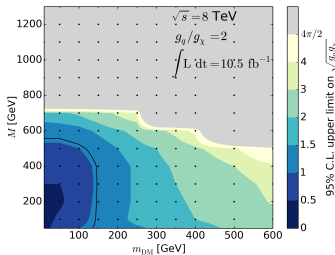
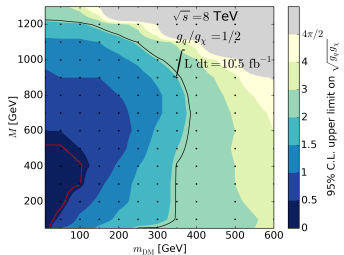
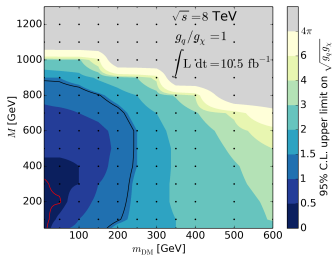
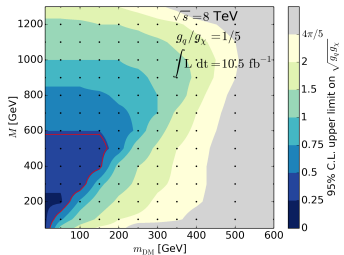


Figure : Example of constraints from 1407.8257.



- ▶ Dark matter is a Big Thing at 14 TeV LHC
- ▶ Need to make sure constraints are robust and can be compared to Direct Detection \Rightarrow EFTs are of limited use
- ▶ Simplified Models give a consistent and robust framework at cost of more parameters

1. Implement Lagrangian in FeynRules
2. Generate parton level events with MadGraph5
3. Match to Pythia 8 for showering
4. Perform detector simulation and analysis in ATOM+Rivet
5. Get out visible cross-section, compare to ATLAS limits

Since we assume axial-vector couplings the minimal width³ is:

$$\Gamma_{\min} = \frac{N_C g_{\text{DM}}^2 M (1 - 4m_{\text{DM}}^2/M^2)^{3/2}}{12\pi} \Theta(M - 2m_{\text{DM}}) + \sum_q \frac{N_C g_q^2 M (1 - 4m_q^2/M^2)^{3/2}}{12\pi} \Theta(M - 2m_q)$$

³Assuming no additional invisible decays.

